

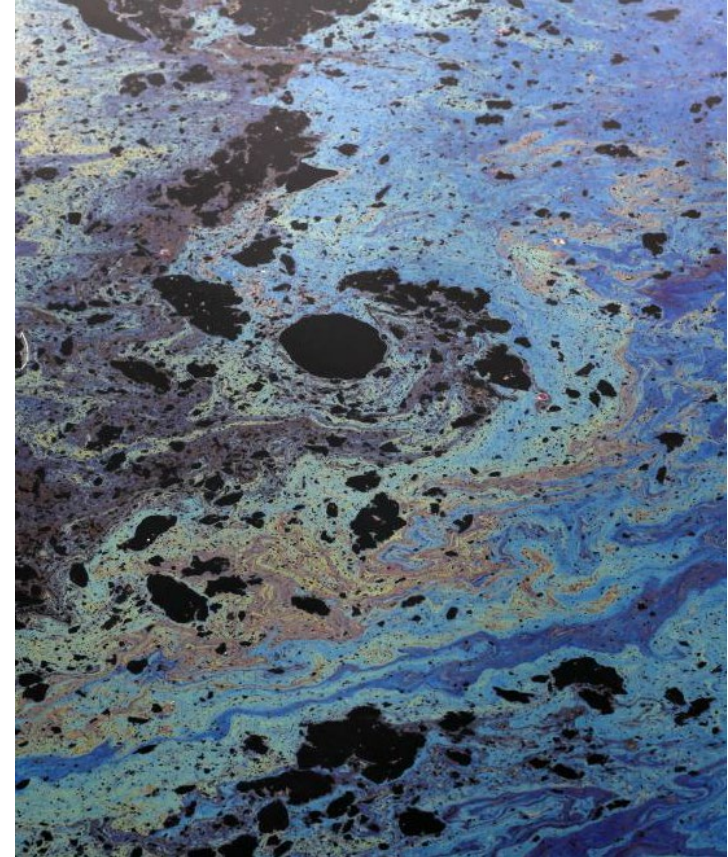
# Effect of Marangoni Forces on Interfacial Heat and Mass Transfer

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# Background

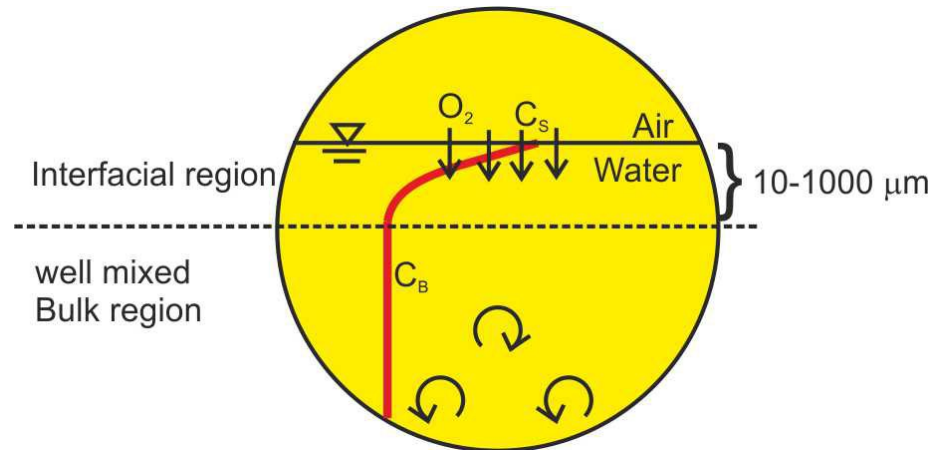
# Interfacial Mass Transfer

Gas transfer ; **molecular diffusion** ↔ **turbulence in the water phase**

**Advective-diffusive** :  $\langle j_z \rangle = - \left[ D \frac{\partial \langle c \rangle}{\partial z} - \langle w' c' \rangle \right]$

j: gas flux  
D: molecular diffusion  
c: concentration  
w: vertical velocity

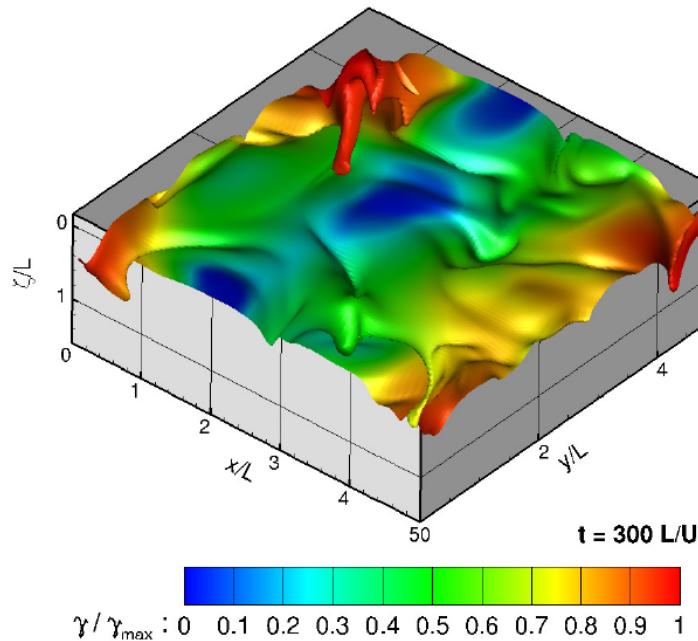
**Gas transfer of low-diffusive gases ( $O_2$ ,  $CO_2$ ) is marked by a very thin concentration boundary layer at the water side**



# Surface contamination

Focus is on interfacial pollution by (insoluble) surfactants

- Surfactants reduce the surface stress of water
- Upwelling and downwelling motions typically lead to non-uniform surfactant concentrations and non-uniform surface stresses
- Resulting in Marangoni forces that counteract surface divergence.



Isosurface at 50%  $C_{\text{sat}}$  coloured by the surface divergence

# Aim

To determine a parametrization of the effect of surface contamination on the transfer function  $K_L$

For a clean (no pollution) interface  $K_L$  scales as

$$K_L \propto Sc^{-1/2}$$

where  $Sc$  is the Schmidt number.

For a very dirty interface

$$K_L \propto Sc^{-2/3}$$

What happens at (very) moderate levels of pollution?

$$K_L \propto Sc^{-q}$$

The power  $q$  will likely depend on  $\frac{Ma}{Ca}$

# Modelling Pollution Effects

Surface tension,  $\sigma$ , depends on the pollutant concentration,  $\gamma$ .

$$\sigma = \sigma(\gamma)$$

After normalization define the Marangoni number by

$$Ma = - \frac{d\sigma}{d\gamma}$$

which we assume to be constant. From the model presented in Shen *et al.*, (2004) JFM, Vol. 506, after some algebra, we obtain:

$$\left. \frac{\partial u}{\partial z} \right|_{interface} = - \frac{Ma}{Ca} \frac{\partial \gamma}{\partial x}$$

$$\left. \frac{\partial v}{\partial z} \right|_{interface} = - \frac{Ma}{Ca} \frac{\partial \gamma}{\partial y}$$

$u$ : x-velocity

$v$ : y-velocity

$Re$ : Reynolds number

$Ca$ : Capillary number

$\gamma$ : surfactant concentr.

# Problem Investigated

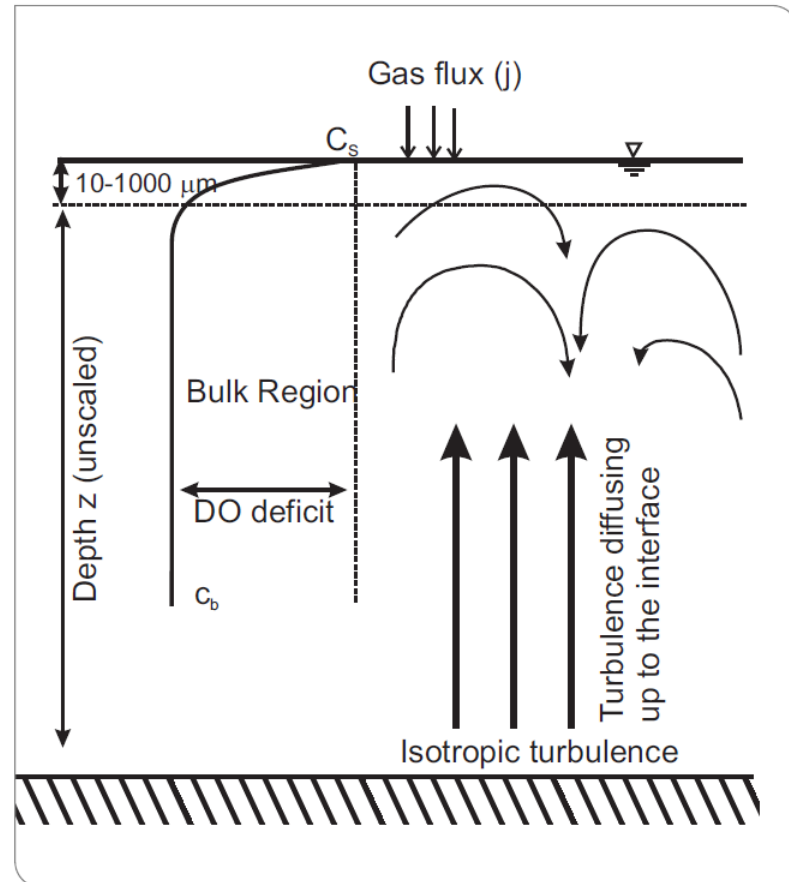
# Physical Problem

## Grid-stirred-driven gas transfer

Convenient analogy to bottom shear induced turbulence



[www.xs4all/rdemming/travel/Indonesia](http://www.xs4all/rdemming/travel/Indonesia)





# Computational Setup

## Boundary conditions

$$\text{Top: } \frac{\partial u_i}{\partial z} \Big|_{\text{top}} = - \frac{Ma}{Ca} \frac{\partial \gamma}{\partial x_i}, i = 1, 2$$

various levels of contaminations

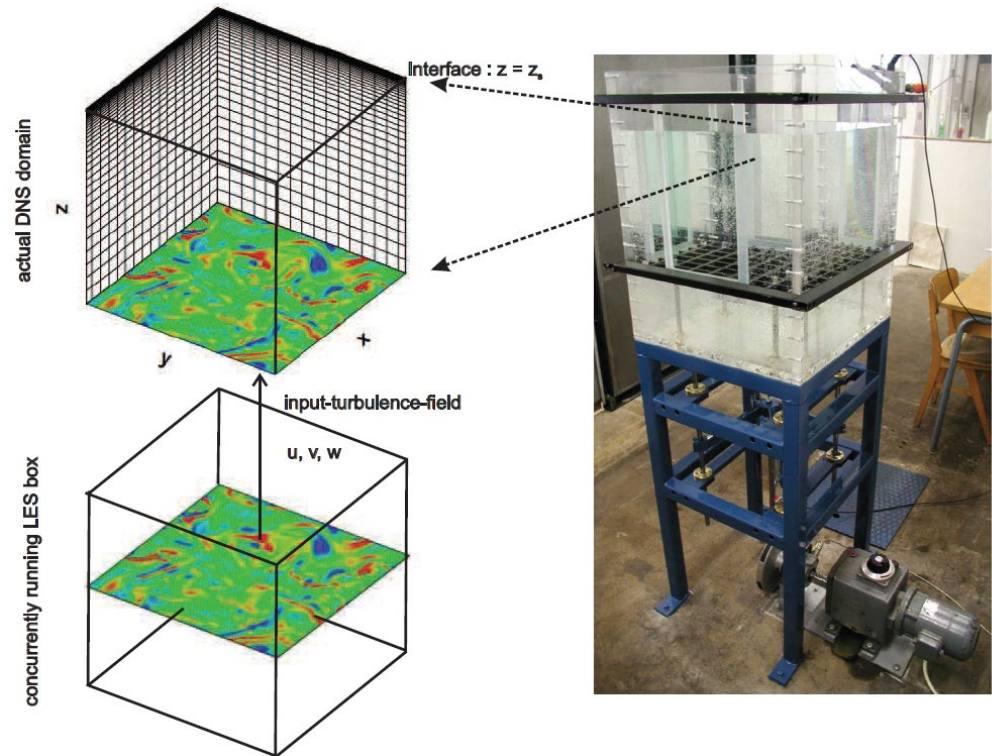
Sides: periodic

Bottom: flow-field copied from LES of isotropic turbulence

$c_{\text{top}} = 1$   
(saturated at all times)

$$c_{\text{bottom}} \rightarrow \frac{\partial c}{\partial z} = 0$$

$x_1, u_1$ :  $x$ ,  $x$ -velocity  
 $x_2, u_2$ :  $y$ ,  $y$ -velocity



# Simulations performed

Case	$U_\infty$	$L_\infty$	$R_T$	$Ma/Ca_T$
S0	0.113	1.033	141	0
S1	0.112	0.958	128	1
S2	0.117	0.994	139	5
S3	0.109	0.984	131	11
S4	0.110	0.927	125	54
S5	0.111	1.021	138	269
SN	0.107	0.898	117	no-slip

$Ma$  is Marangoni number;  $Ca_T = \mu U_\infty / \sigma$  is turb. capillary number

All simulations: 128 x 128 x 212 mesh for 5L x 5L x 3L box

Mesh is refined in z-direction towards surface

Schmidt numbers  $Sc = 2 \dots 32$ ; Surfactant:  $Sc = 2$

Turbulent flow with  $Tu = 40\%$  introduced at bottom

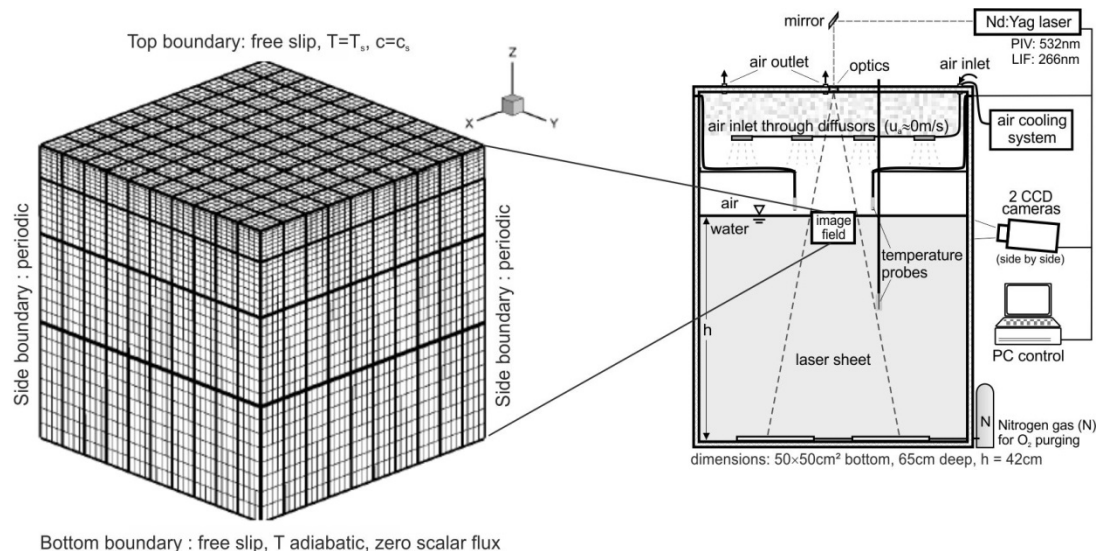
# Numerical Method

Flow fields in main DNS and LES isobox are solved using fourth-order discretisations of convection and diffusion.

A dual mesh strategy is used where up to five scalars can be solved simultaneously on a refined mesh

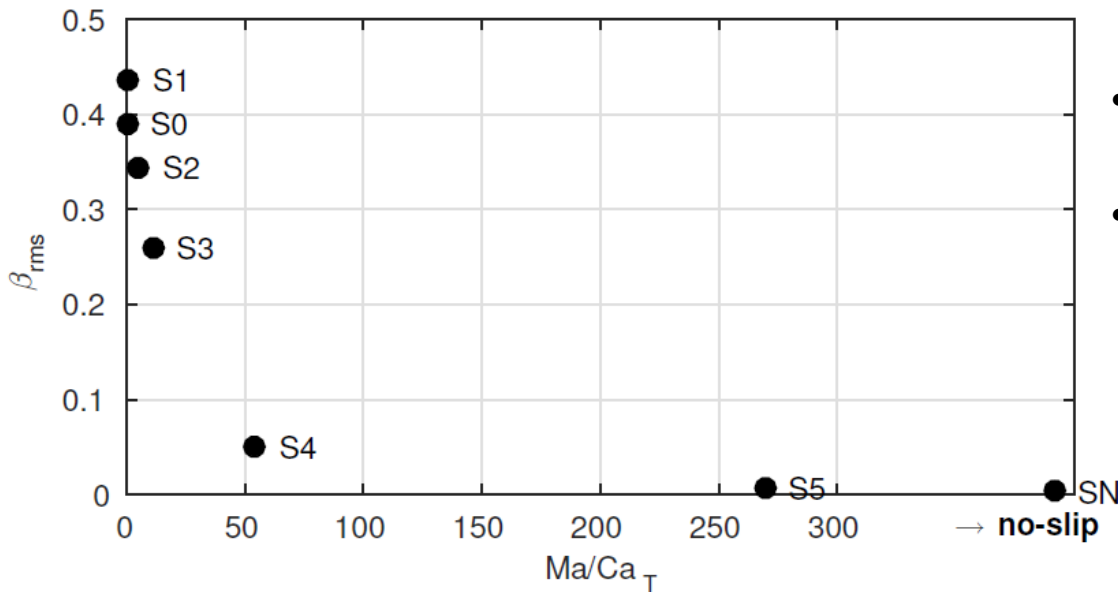
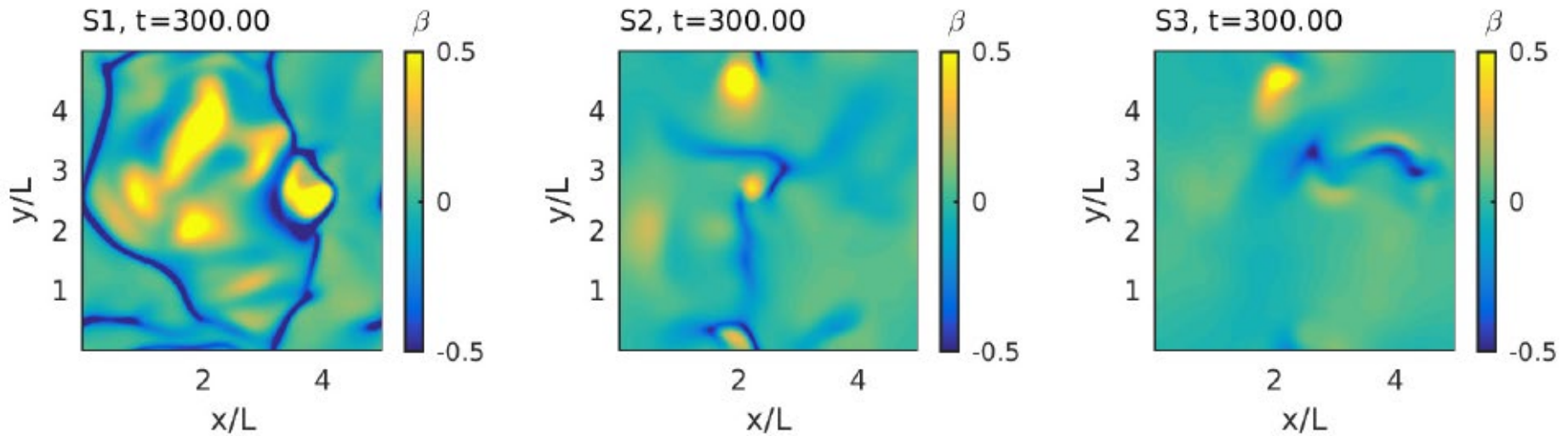
A fifth-order-accurate WENO scheme is used for scalar convection, combined with a fourth order central discretisation for scalar diffusion (same in 2D for surfactant).

Standard Message passing interface (MPI) is applied for communication between blocks.



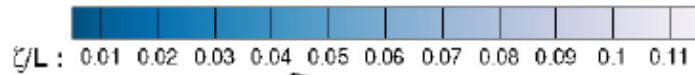
# Effect on near-surface hydrodynamics

# Surface Divergence ( $\beta$ )

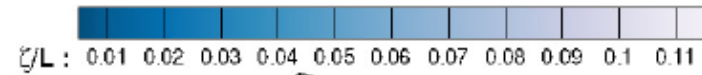
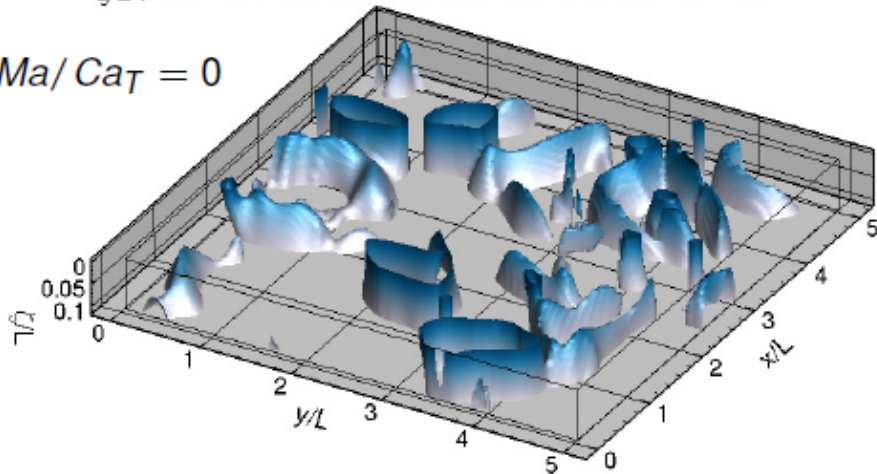


- $\beta$  reduces sharply
- Surface divergence model ( $K_L \propto \sqrt{D\beta_{rms}}$ ) only works for a small range of  $Ma/Ca_T$

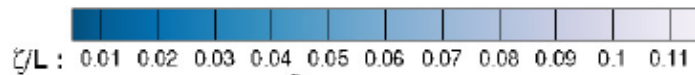
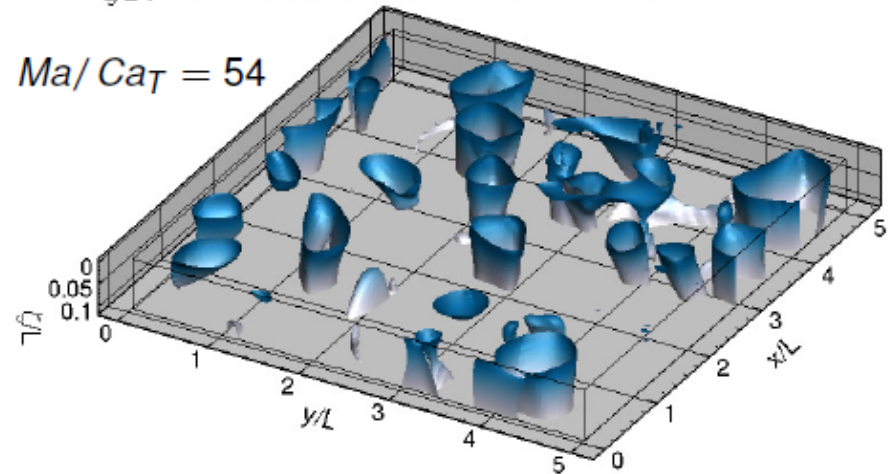
# Instantaneous shear



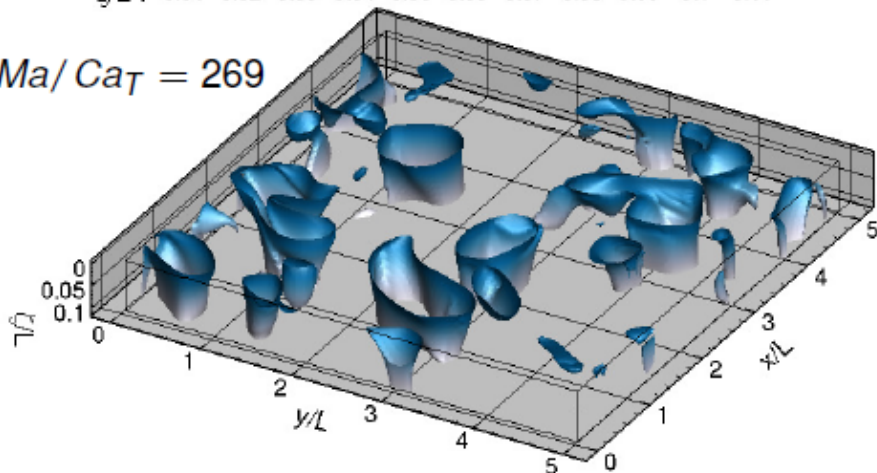
$Ma/Ca_T = 0$



$Ma/Ca_T = 54$



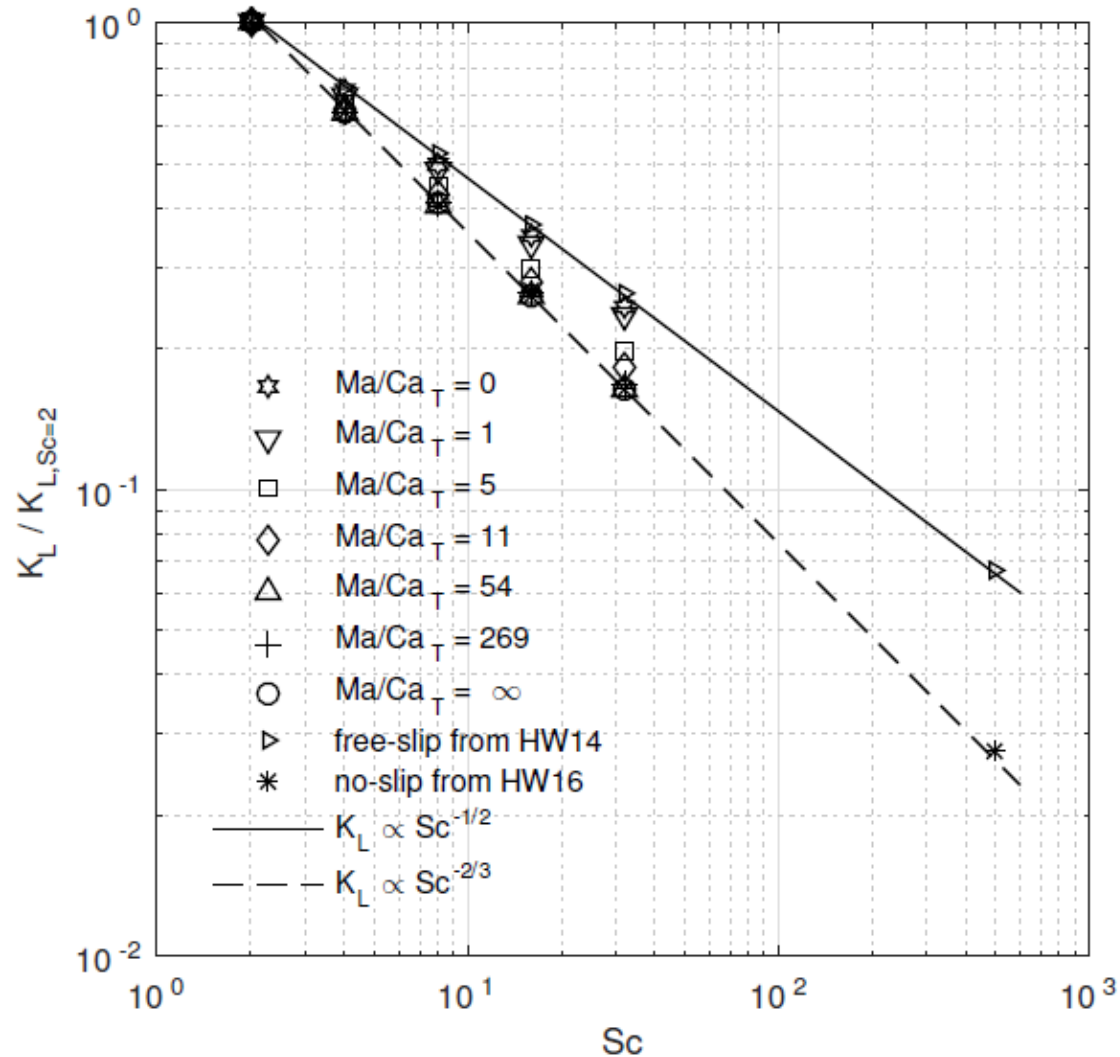
$Ma/Ca_T = 269$



- orientation and cross-sectional area become different
- explain the apparent increase in the integral length scale
- strong correlation between low concentration regions and strong positive surface divergence is lost

# Consequences on interfacial mass transfer

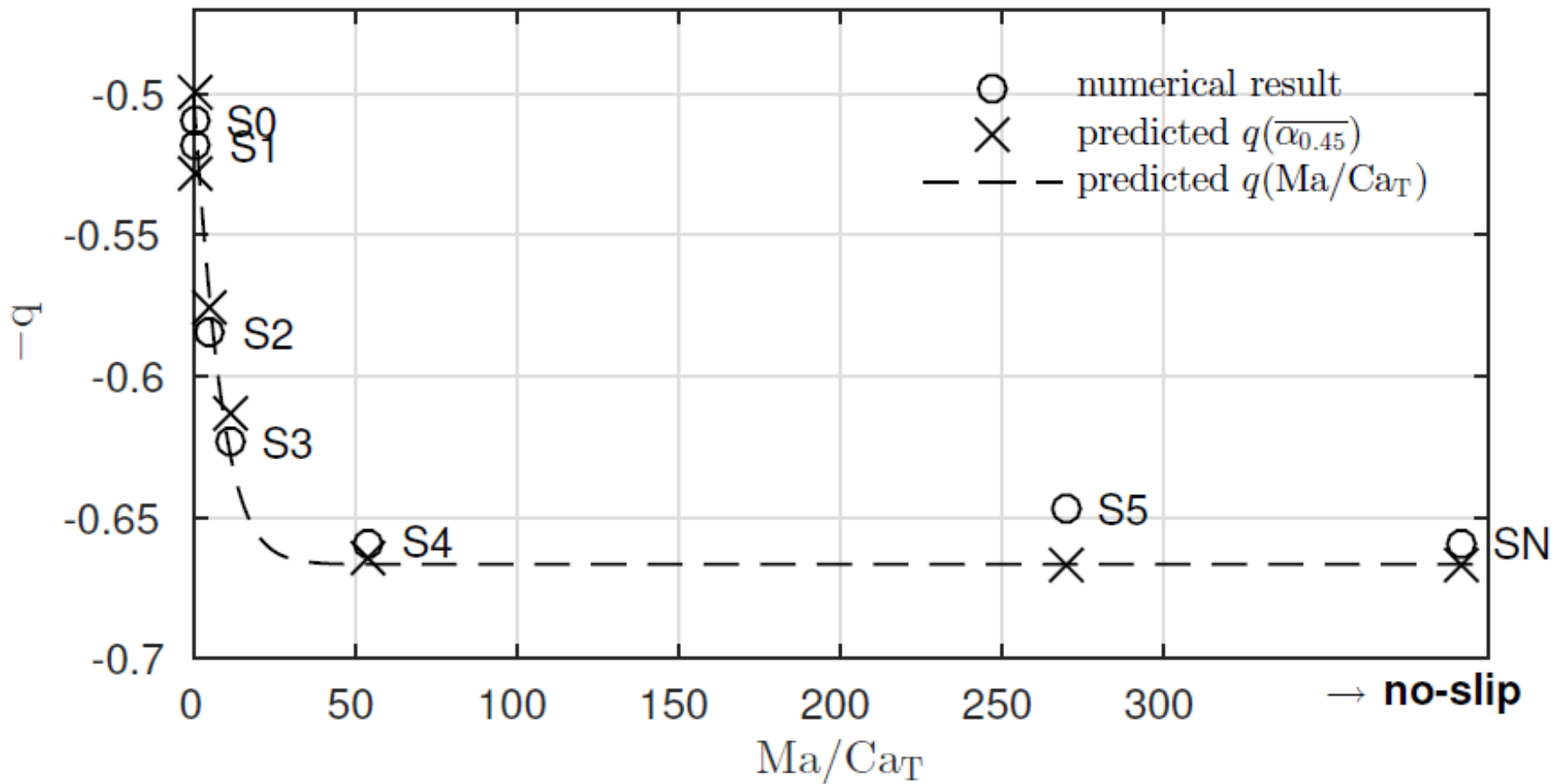
# $K_L$ reduction



- Reduction of  $K_L$
- For each  $Ma/Ca_T$  :  
 $K_L \propto Sc^{-q}$

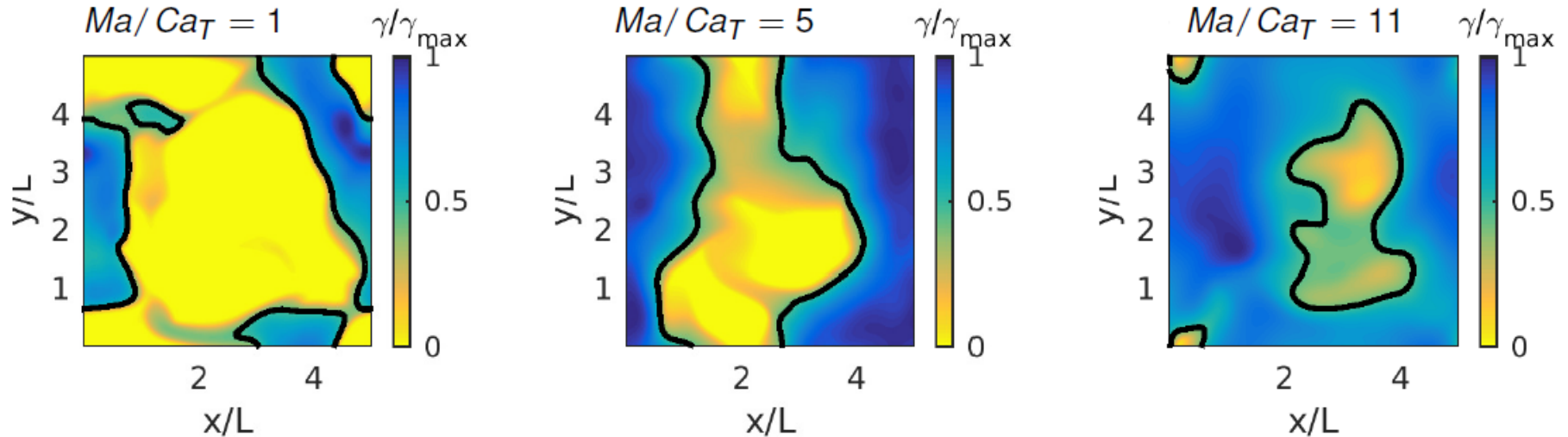


# Transition of $q$



Steep transition of  $q$  from  $\frac{1}{2}$  (clean) to  $\frac{2}{3}$  (very dirty)

# Clean surface fraction



A good correlation was found between the Schmidt exponent  $q$  and clean surface fraction.

# Model

First we assume that for any surface condition

$$K_L = cSc^{-q}R_T^{-r}$$

Clean regions behave as a free-slip boundary ( $q = 1/2$ ), while “dirty” regions behave as a no-slip boundary ( $q = 2/3$ )

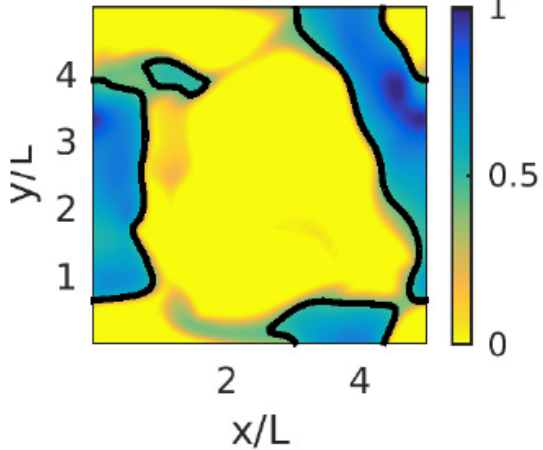
$$cSc^{-q} = \bar{\alpha}c_fSc^{-1/2} + (1 - \bar{\alpha})c_nSc^{-2/3}$$

Use Taylor series expansions to obtain first order appr. of  $c$  and  $q$  that are independent of  $Sc$

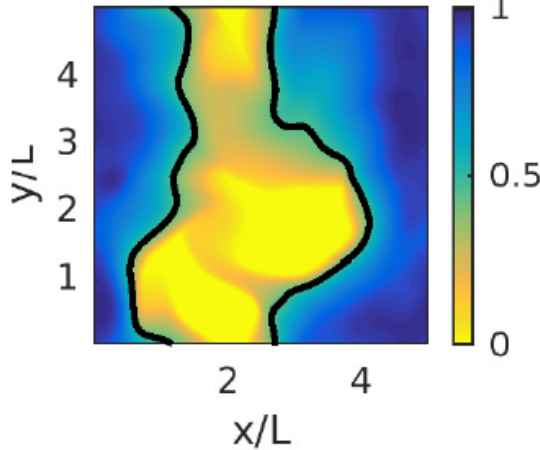
$$c = \bar{\alpha}c_f + (1 - \bar{\alpha})c_n$$
$$q = \frac{2}{3} - \frac{\bar{\alpha}c_f}{6c}$$

# Definition of clean

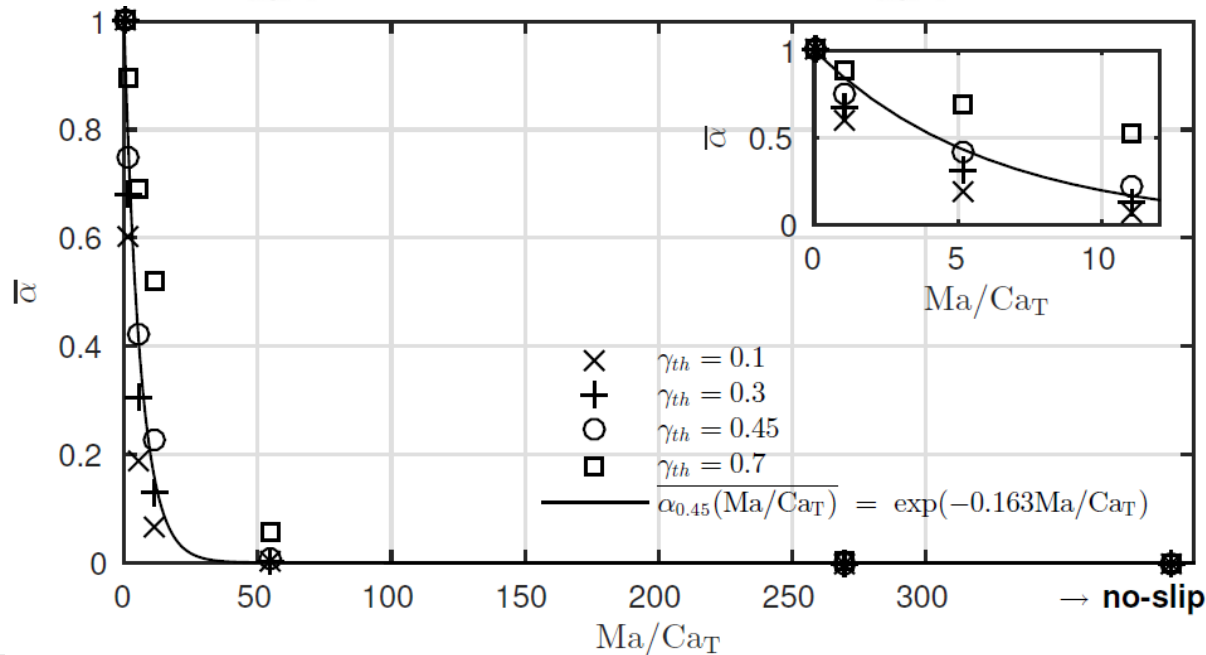
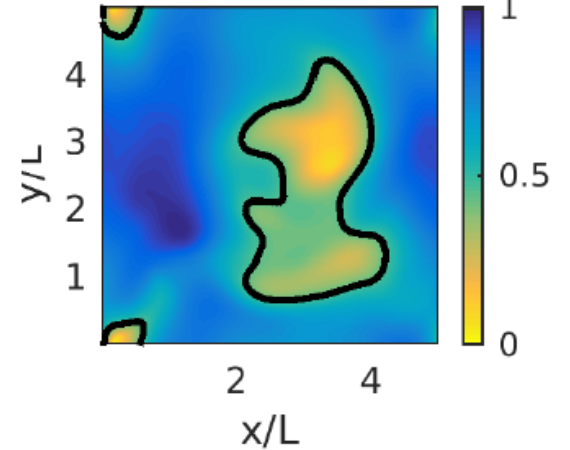
$Ma/Ca_T = 1$   $\gamma/\gamma_{max}$



$Ma/Ca_T = 5$   $\gamma/\gamma_{max}$



$Ma/Ca_T = 11$   $\gamma/\gamma_{max}$

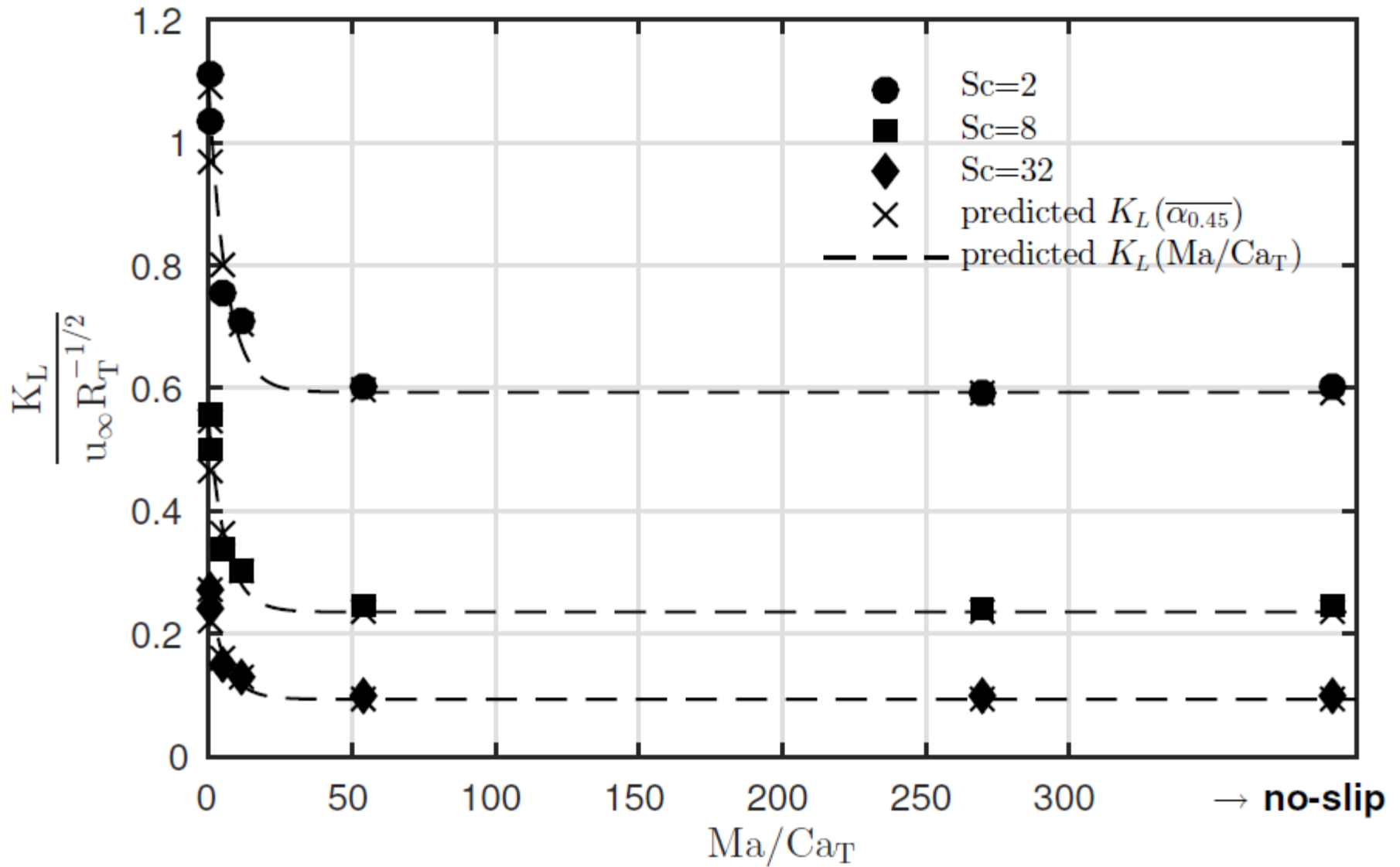


$\gamma$  : surfactant concentr.

$\alpha$  : clean surface fraction

$\gamma_{th} = \gamma/\gamma_{max}$  is threshold

# Evaluation of model



# Conclusions

- It was confirmed that the even small levels of surfactant contamination have a large effect on heat and gas transfer
- With increasing  $Ma/Ca_T$ , the surface divergence,  $\beta$ , becomes progressively damped
- Resulting in a quick transition to a  $K_L \propto Sc^{-2/3}$  scaling which is typical for a no-slip surface
- The transition can be linked to the mean clean surface fraction which is a relatively easily observable parameter.

# Future work

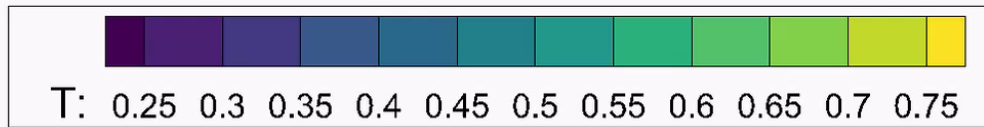
# Marangoni Forces promoting Buoyant Instability

- Surface cooling due to evaporation can be modelled by applying a constant heat flux at the surface.
- A buoyant instability results from this unstable layering of the water as cold water is heavier than warm water.
- The developing buoyant instability results in horizontal gradients in the surface temperature
- As surface tension reduces with increasing temperature, Marangoni forces are generated that act to promote the buoyant instability



# Marangoni Forces promoting Buoyant Instability

Top plane view, heat flux =  $-\frac{2.33}{RePr}$ ,  $Re = 100$ ,  $Pr = 7$



t=10.25

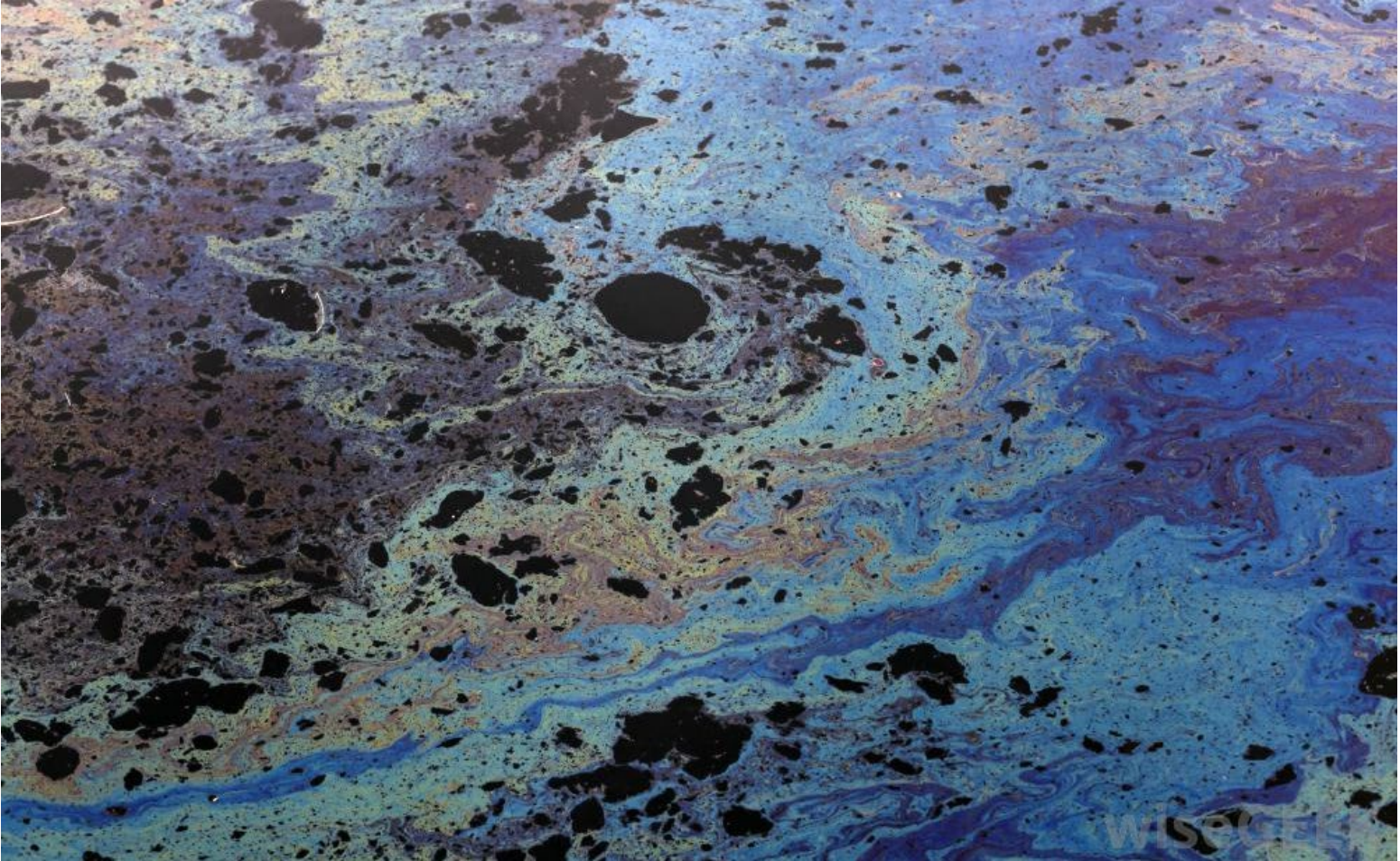
Ma = 0



Ma = 1



# The end



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