

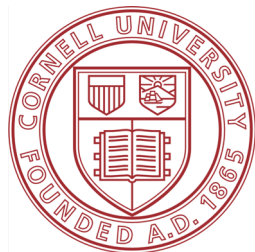
Laboratory investigation of significant gas transfer enhancement via capillary-gravity bow waves

Katherine E. Adler

PhD Candidate at Cornell University

Edwin A. Cowen

Professor at Cornell University



Using standard definition of k and δ

$$F = k\Delta C = k\left(\frac{C_a}{H} - C_b\right) = k(C_s - C_b)$$

F = mass flux/unit surface area ($M L^{-2} T^{-1}$)

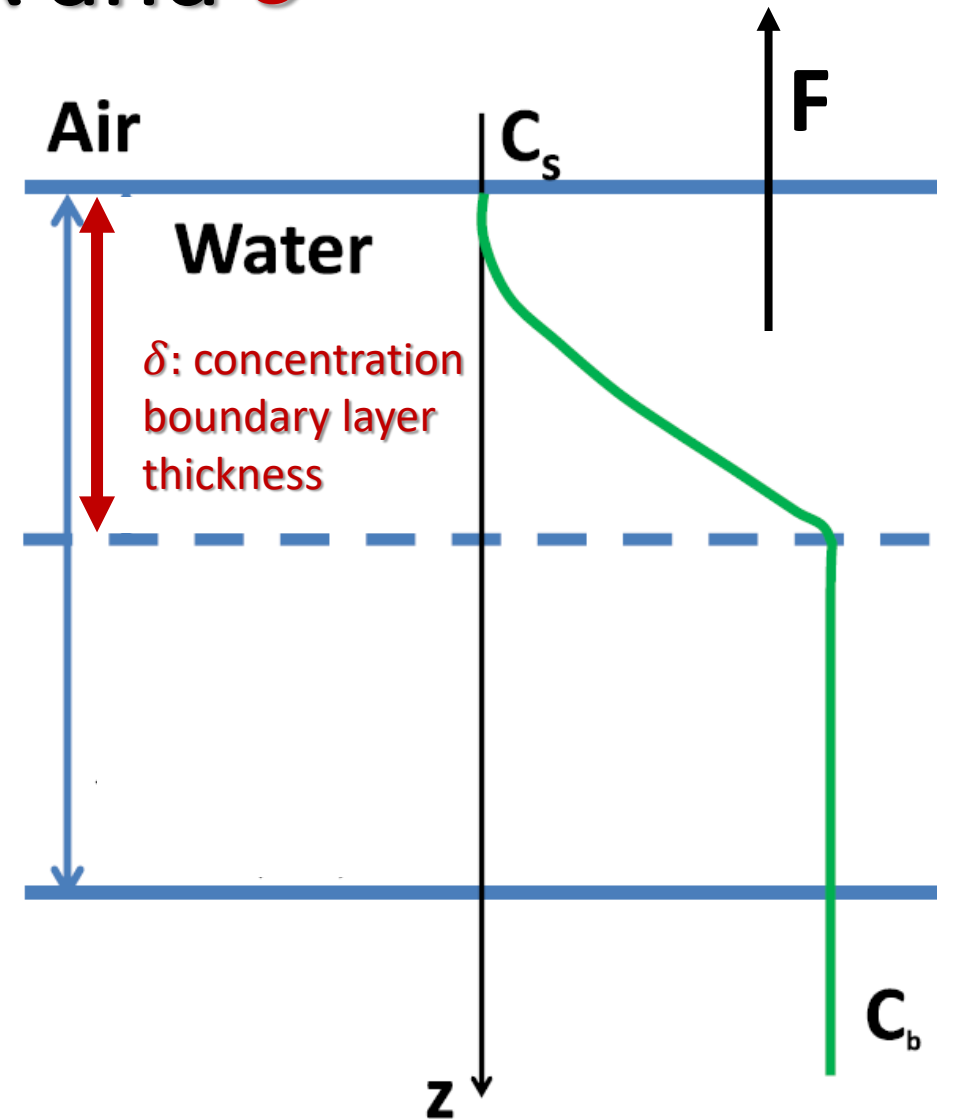
k = transfer velocity ($L T^{-1}$)

ΔC = concentration gradient over CBL ($M L^{-4}$)

C_a = the bulk air concentration ($M L^{-3}$)

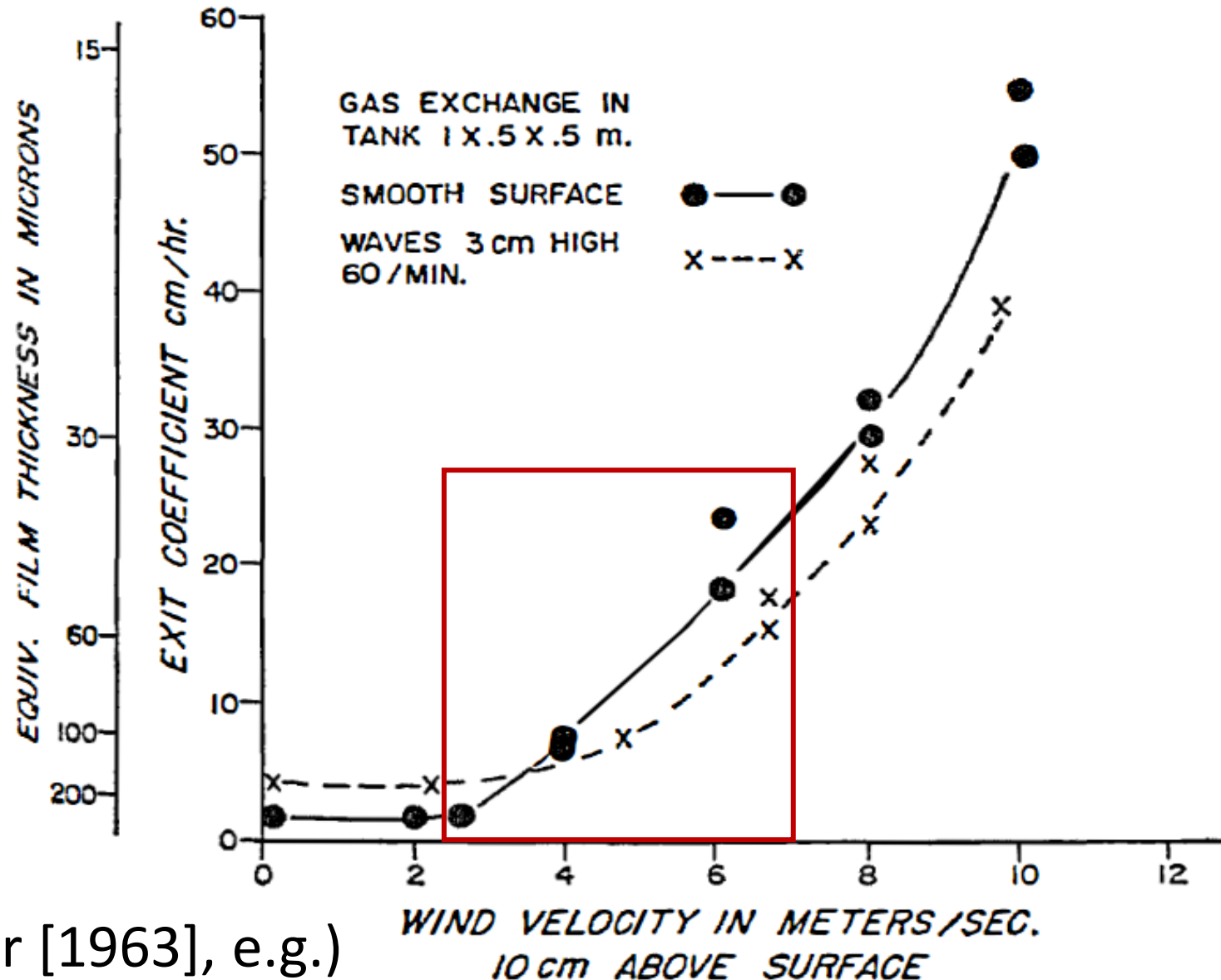
C_b = the bulk water concentration ($M L^{-3}$)

H = Henry's Law constant



k is enhanced at onset of wind waves

For $U_{10} < 7 \frac{m}{s}$
capillary waves
dominate
(e.g., Bourasa *et al.*, 1999)



(Figure from Kanwisher [1963], e.g.)

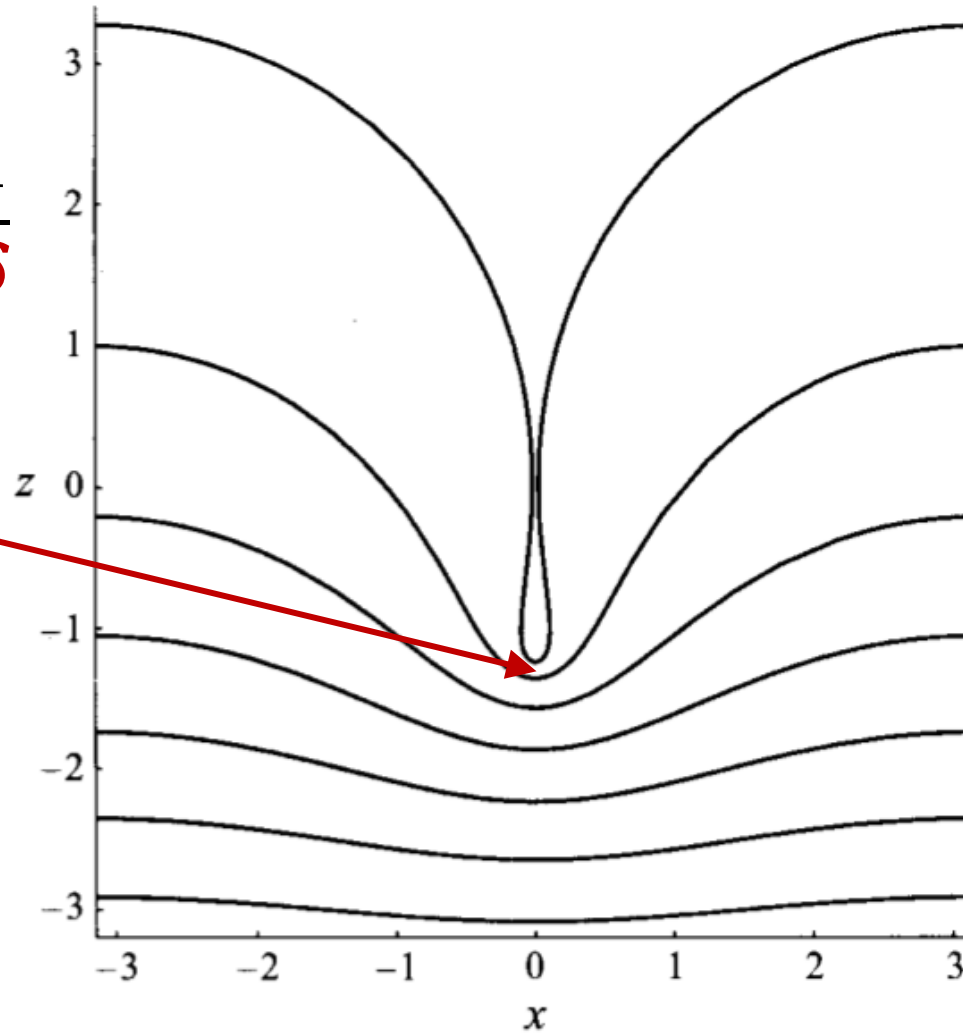
Capillary waves thin δ , increase concentration gradient

Fick's Law:

$$F = D \frac{\partial C}{\partial z} \rightarrow k \propto \frac{1}{\delta}$$

Straining troughs
leads to thinner

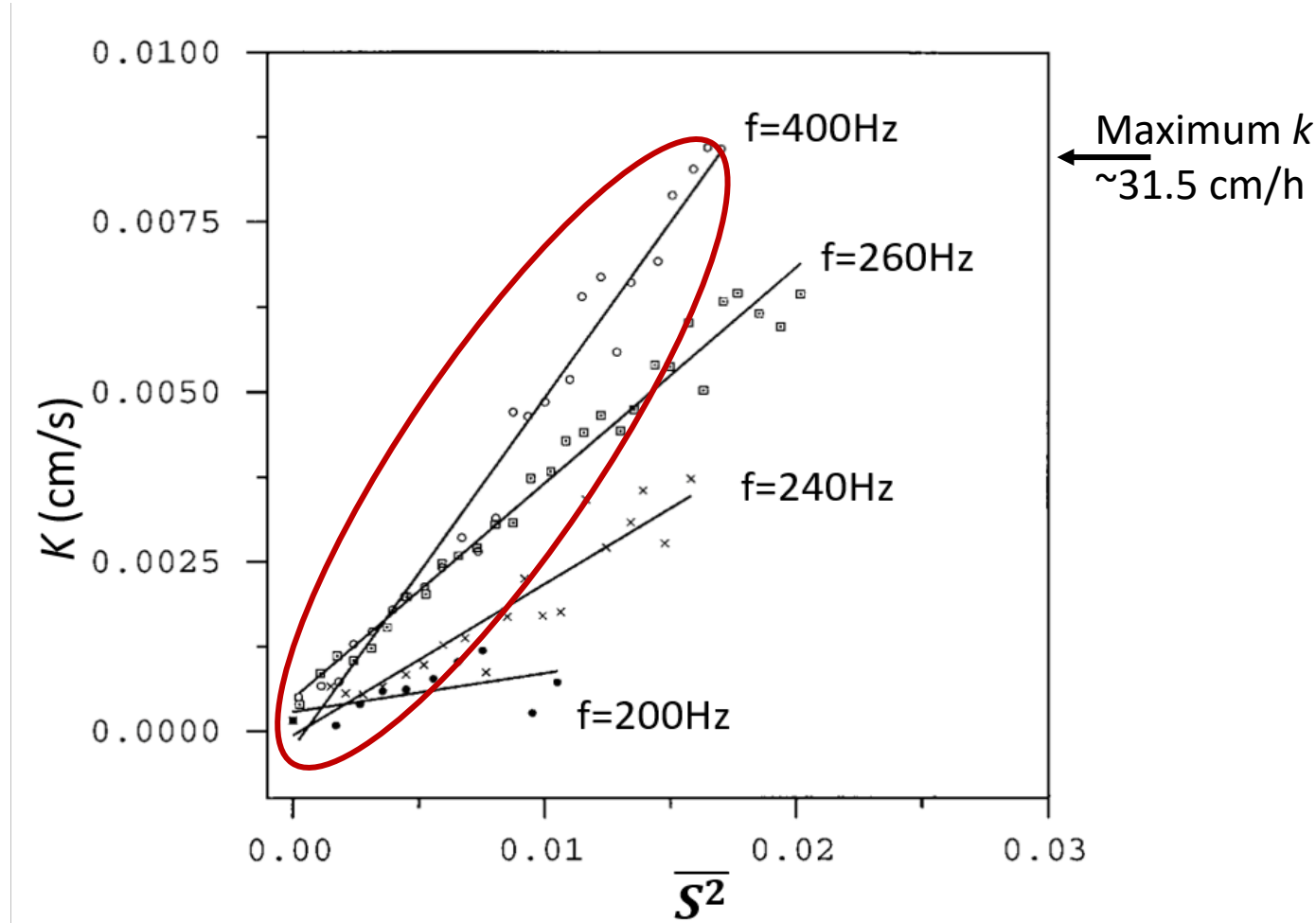
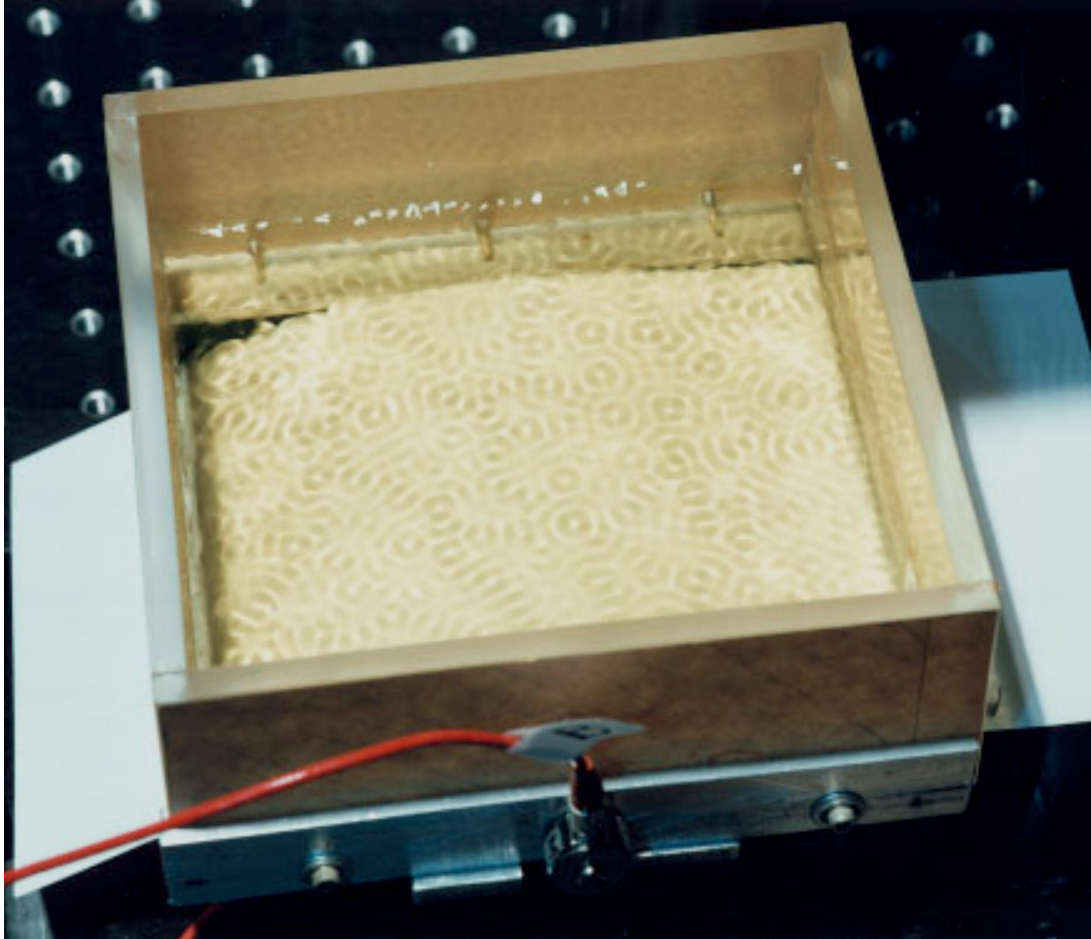
δ



Szeri model estimates
3-fold enhancements

The surface elevation and streamlines of the exact solution for a capillary wave amplitude $A = 0.45$ (Szeri, 1997).

Small-scale model \rightarrow almost 100x increase in k



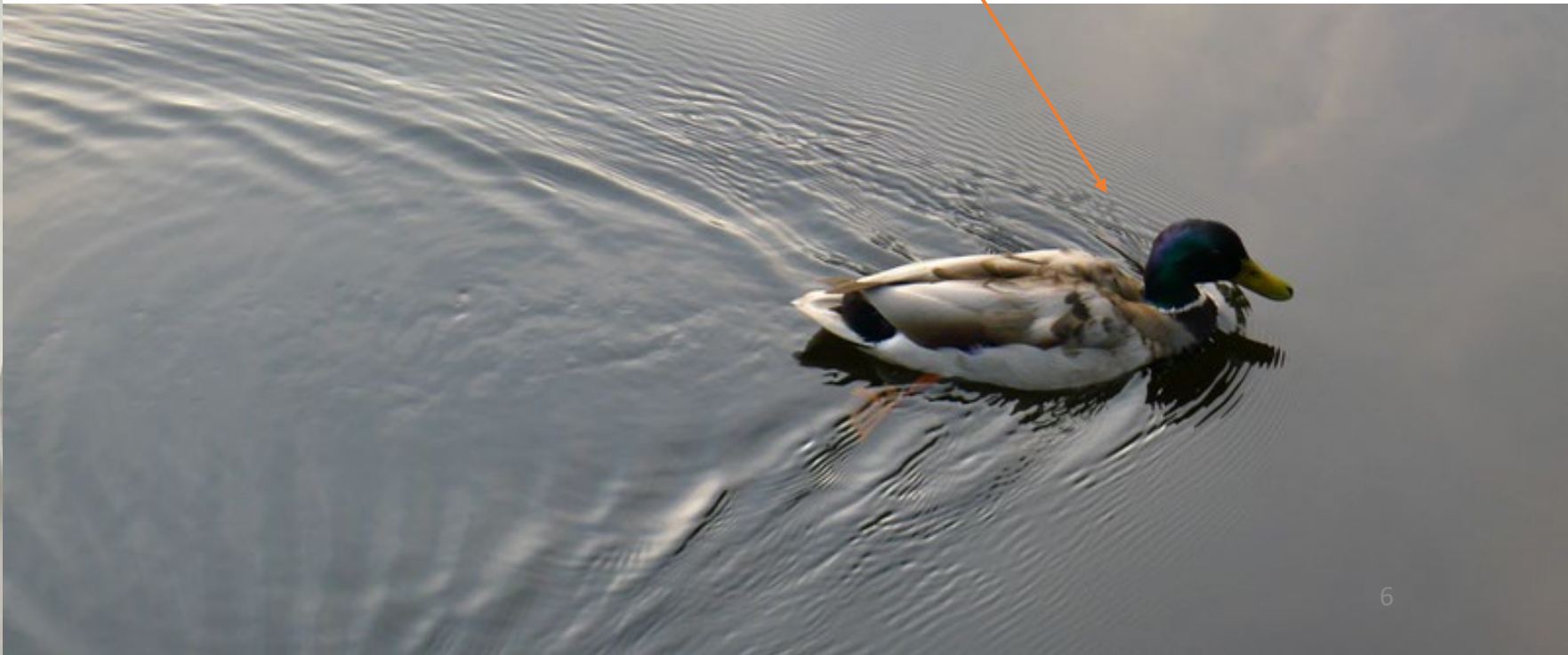
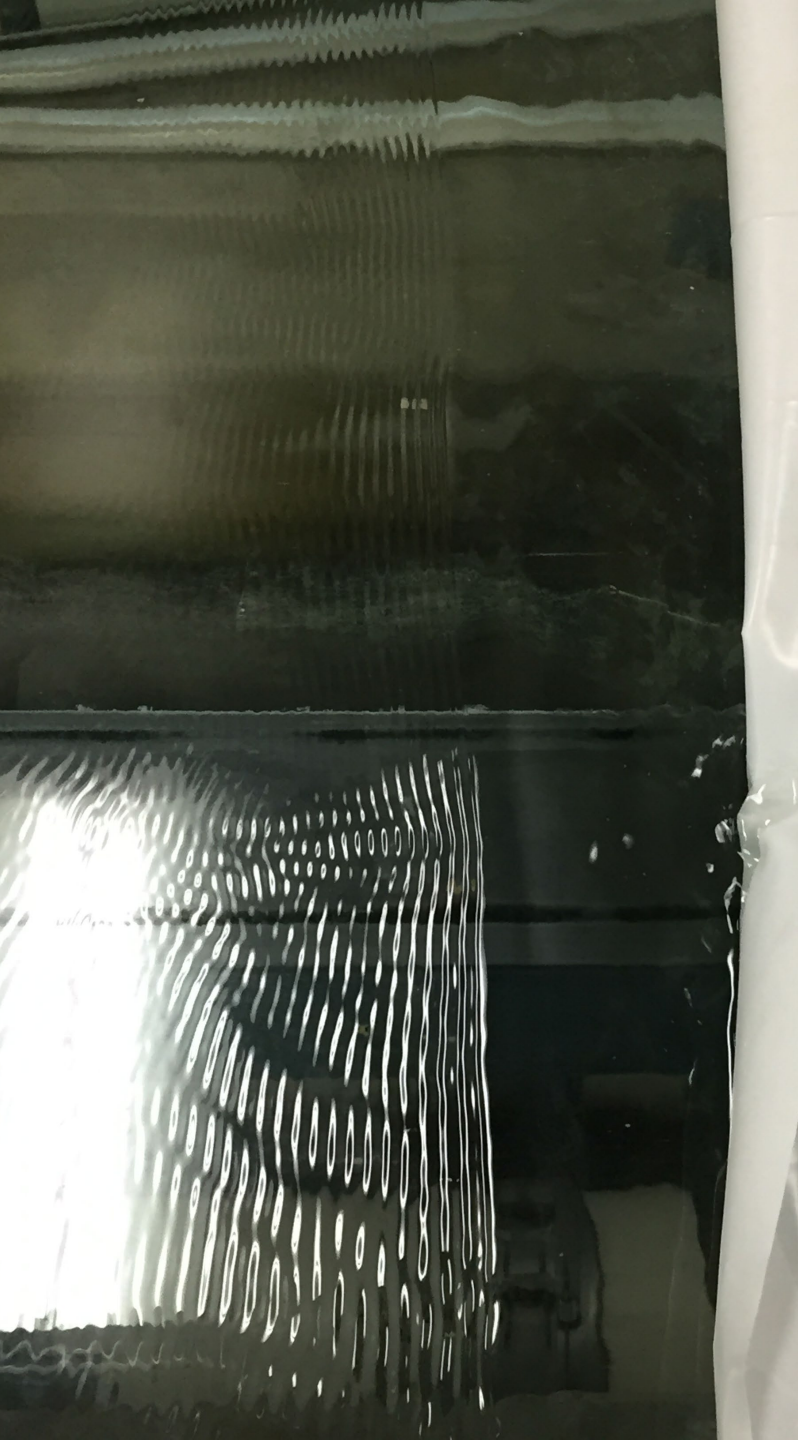
S = the slope of the water surface

(Saylor & Handler, 1999)

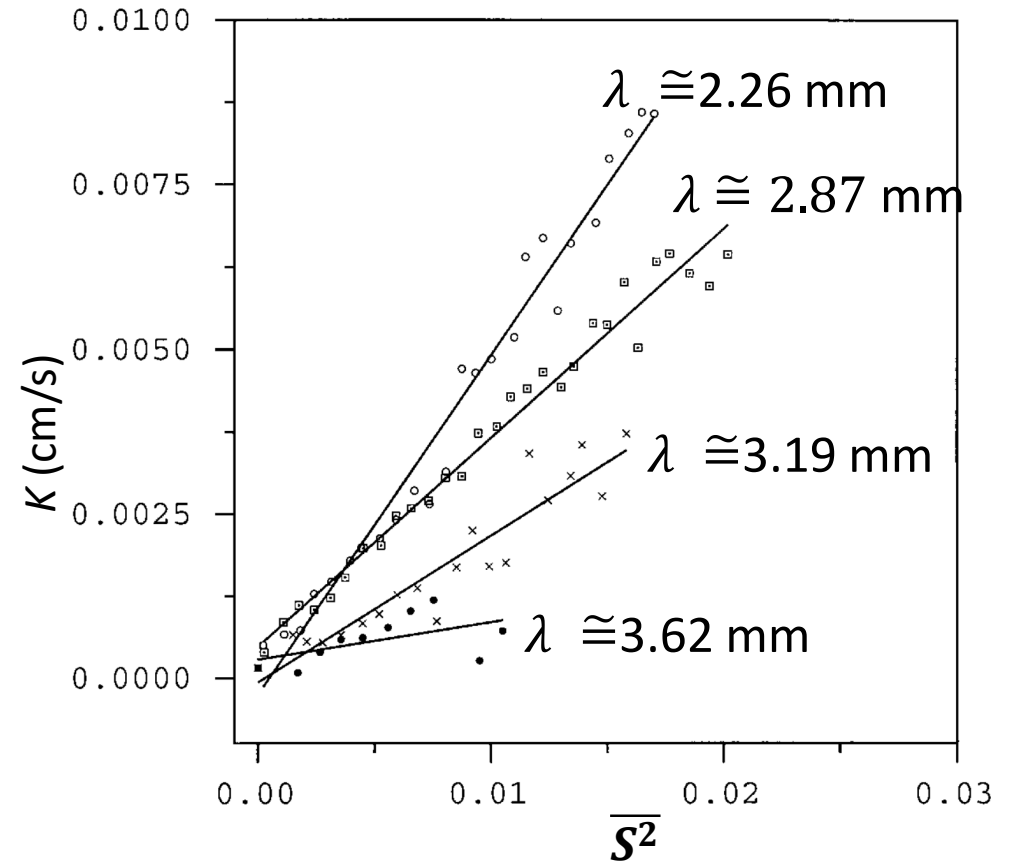
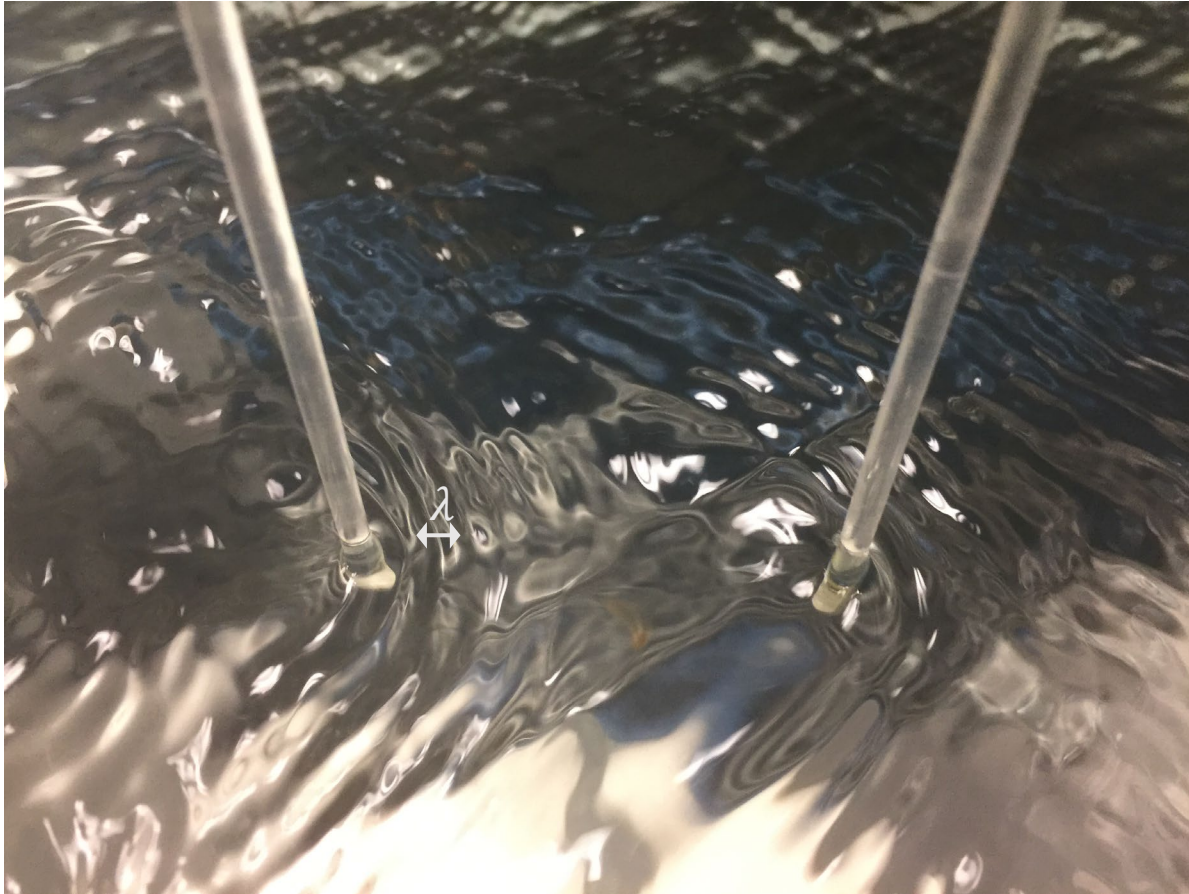
$2.26 \text{ mm} \leq \lambda \leq 3.62 \text{ mm}$

We expand to a larger scale by
generating capillary-gravity bow
waves

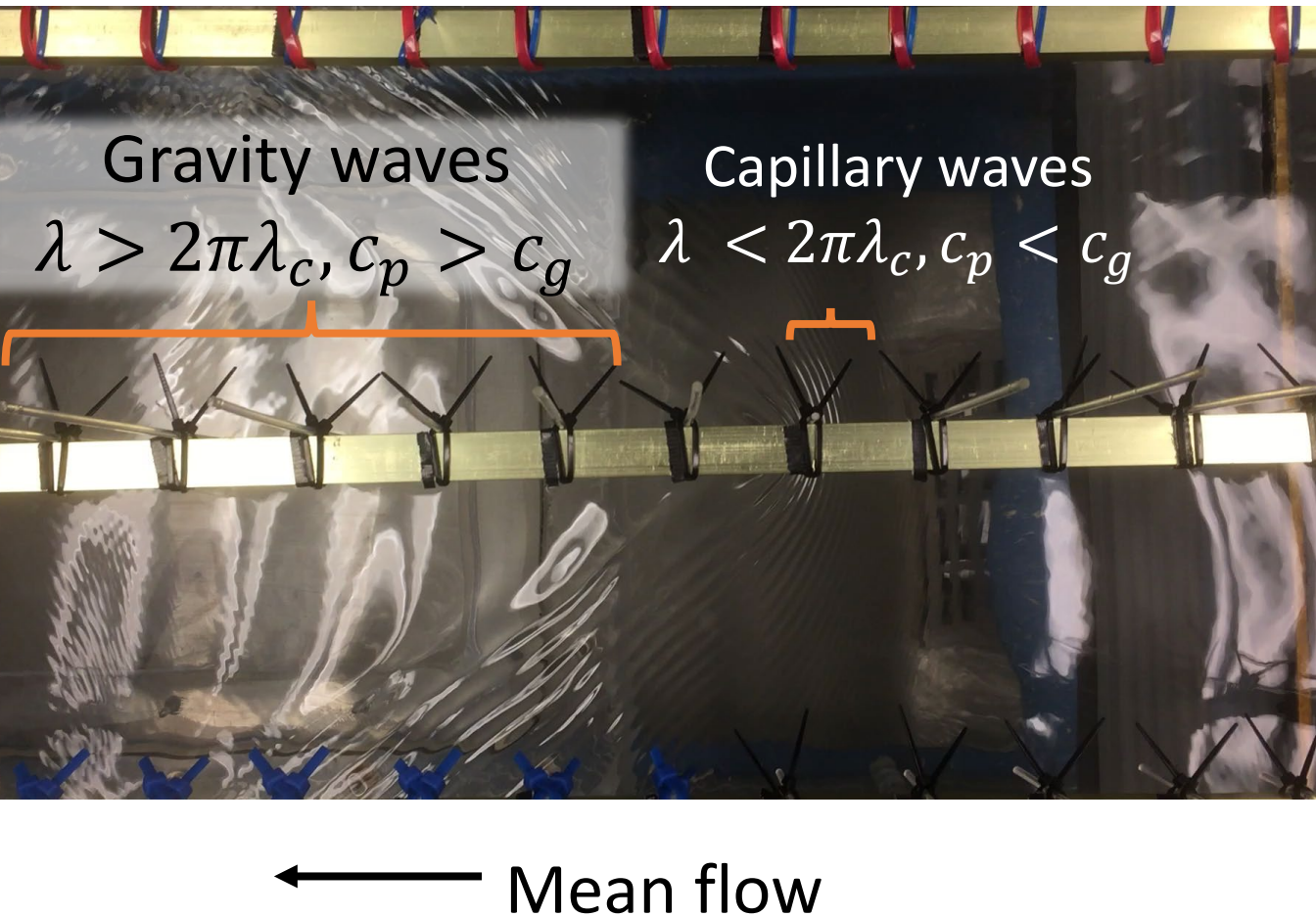
Capillary-gravity
bow waves
(mm-cm-scale)



Greatest k enhancement at shortest wavelength



Analytical model to target relative velocities



$$\lambda_c = \sqrt{\frac{\gamma}{\rho g}} \quad \leftarrow \text{Surface tension}$$

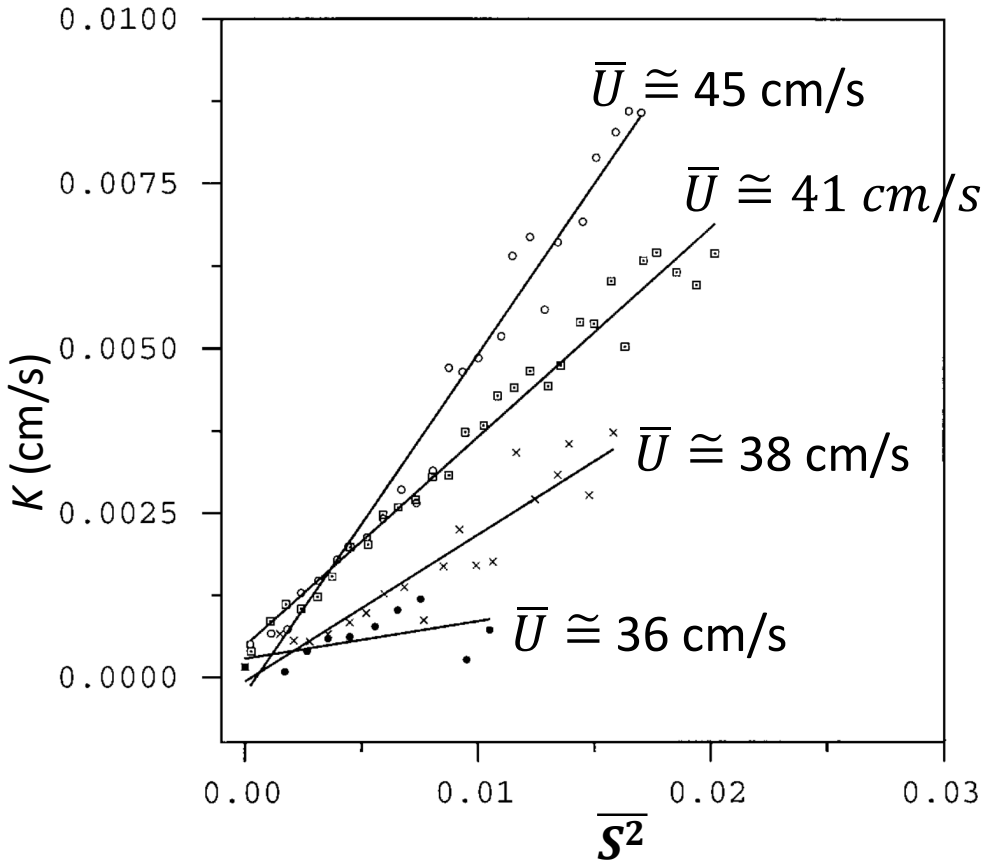
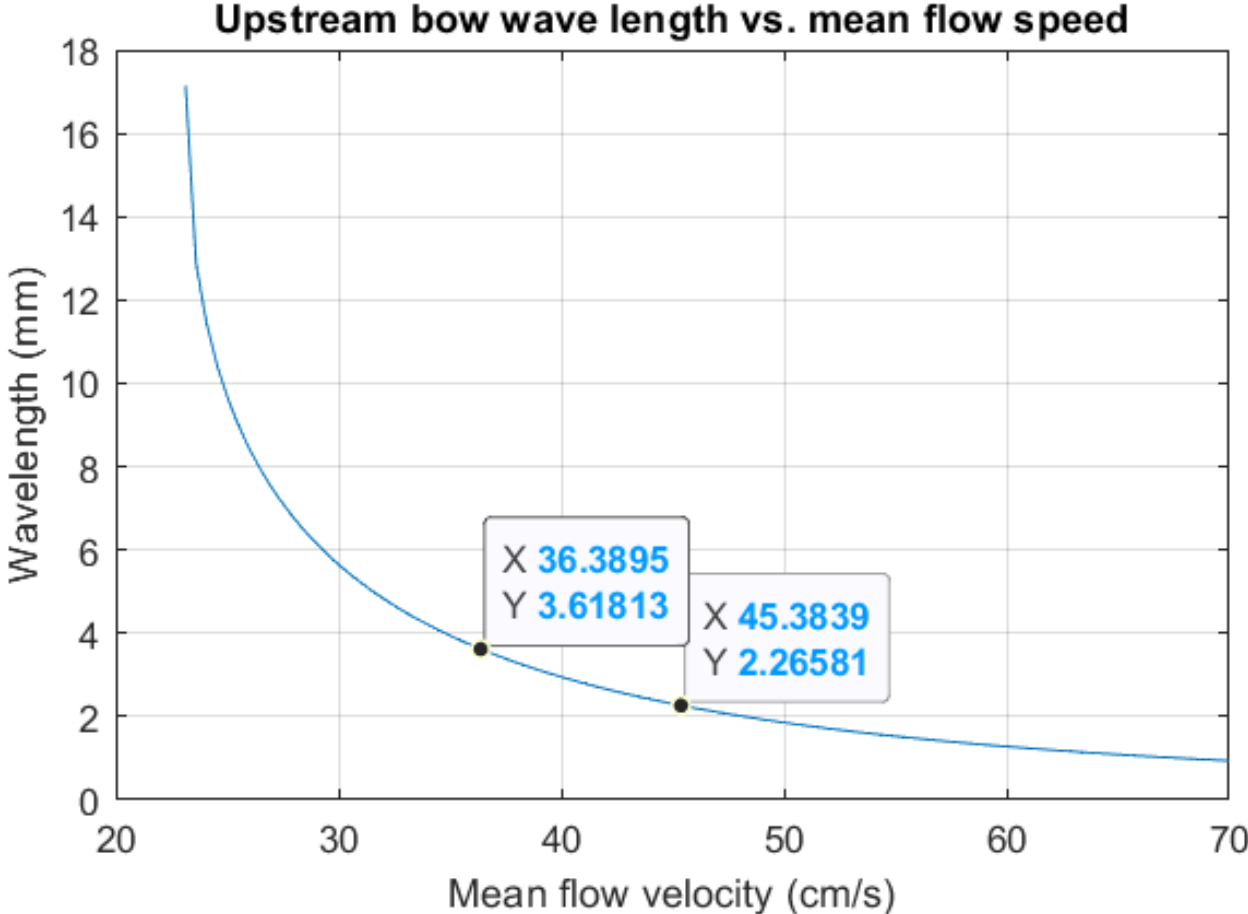
$$c_p = \sqrt{\frac{g}{K} + \frac{\gamma K}{\rho}} \quad \leftarrow \text{Wave number}$$

$$c_g = \frac{g + \frac{3\gamma K^2}{\rho}}{2\sqrt{gK + \frac{\gamma K^3}{\rho}}}$$

$$K_{C/G} = \frac{1}{\lambda_c} U^2 \sqrt{\frac{\rho}{4g\gamma}} \left\{ 1 \pm \sqrt{1 - \frac{4g\gamma}{\rho U^4}} \right\}$$

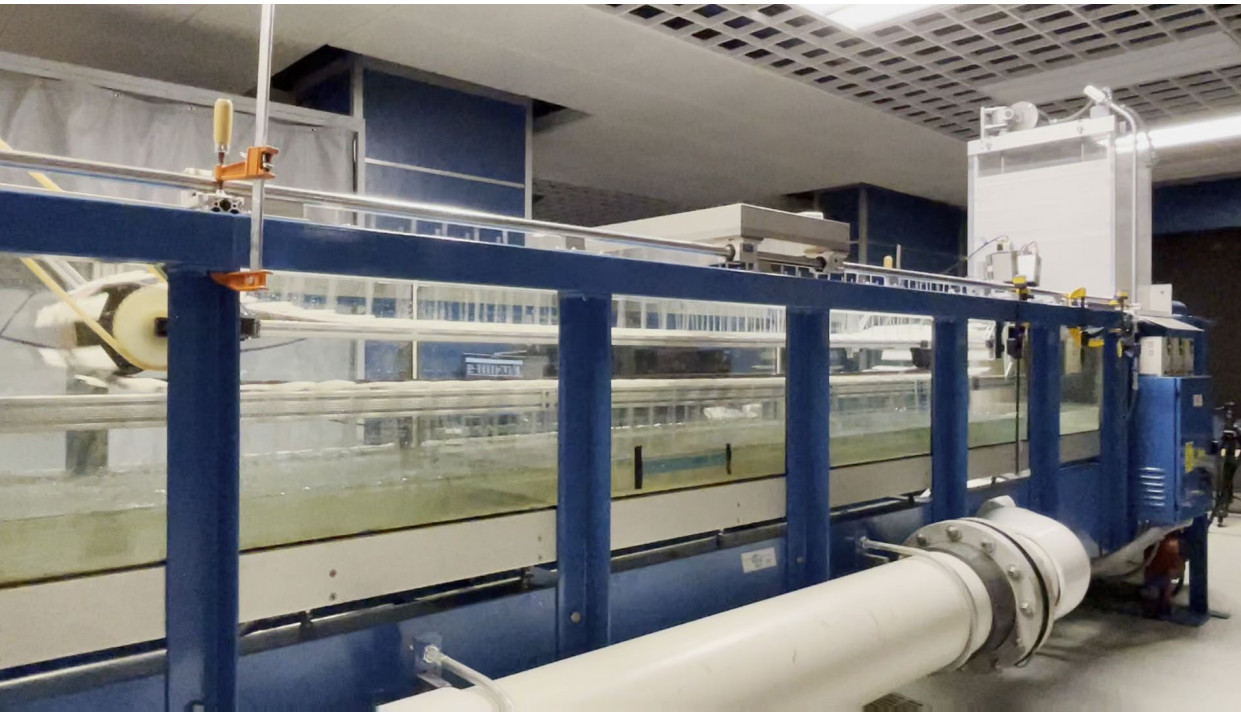
(Raphaël and De Gennes, 1995)

Target velocities around 36-45 cm/s and higher

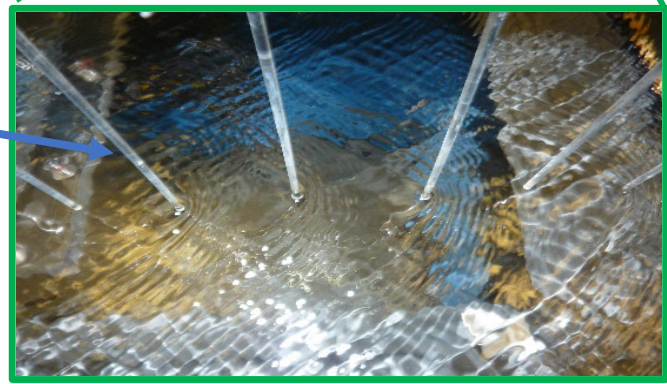
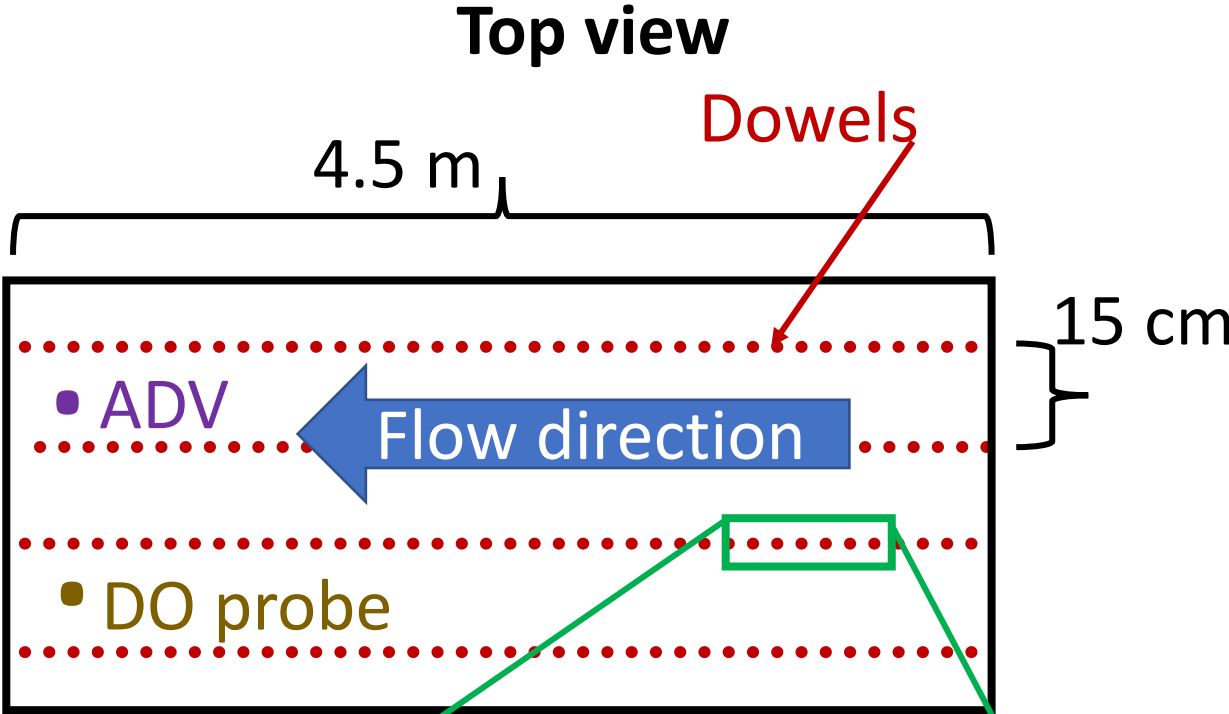


(Right figure modified from Saylor & Handler, 1999)

Dowel array to populate entire surface



Recirculating flume with suspended dowel array



3.2-mm dowels

Remove dissolved oxygen and measure reaeration

Steps:

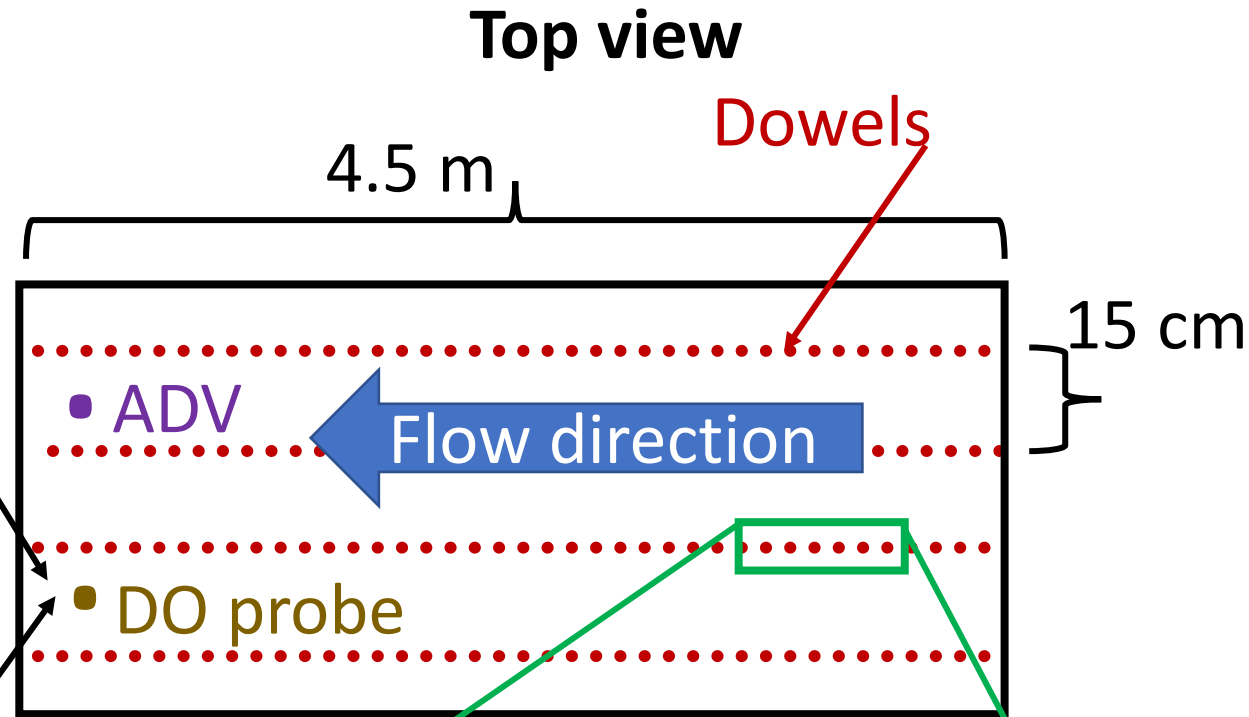
- Use $CoCl_2$ and Na_2SO_3 to remove oxygen
- Measure voltage, temperature, pressure over time
- Use temperature-sensitive calibration to convert voltage to concentration
- Solve for theoretical saturation concentration



YSI DO Probe



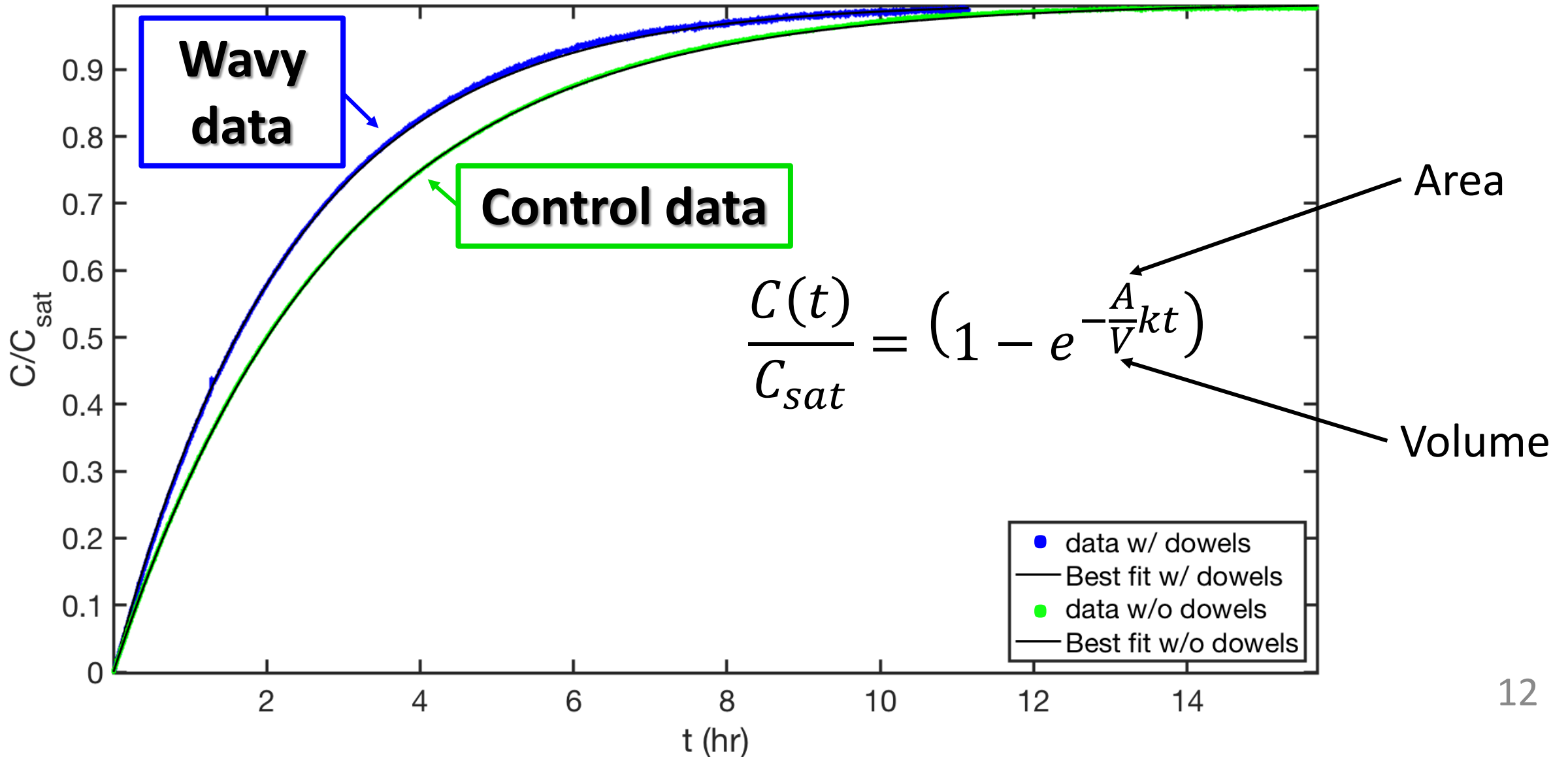
PME miniDOT



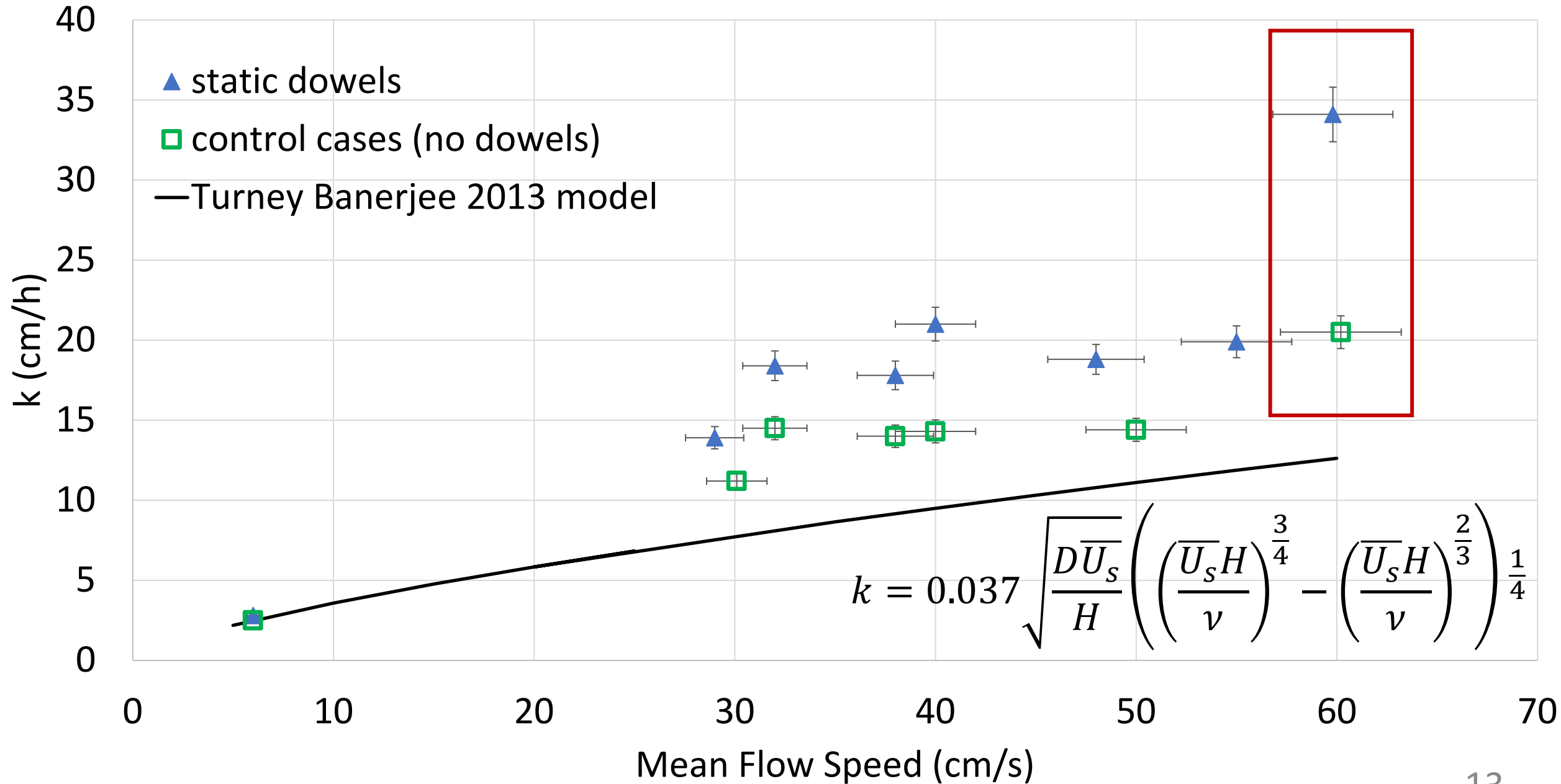
3.2-mm
dowels



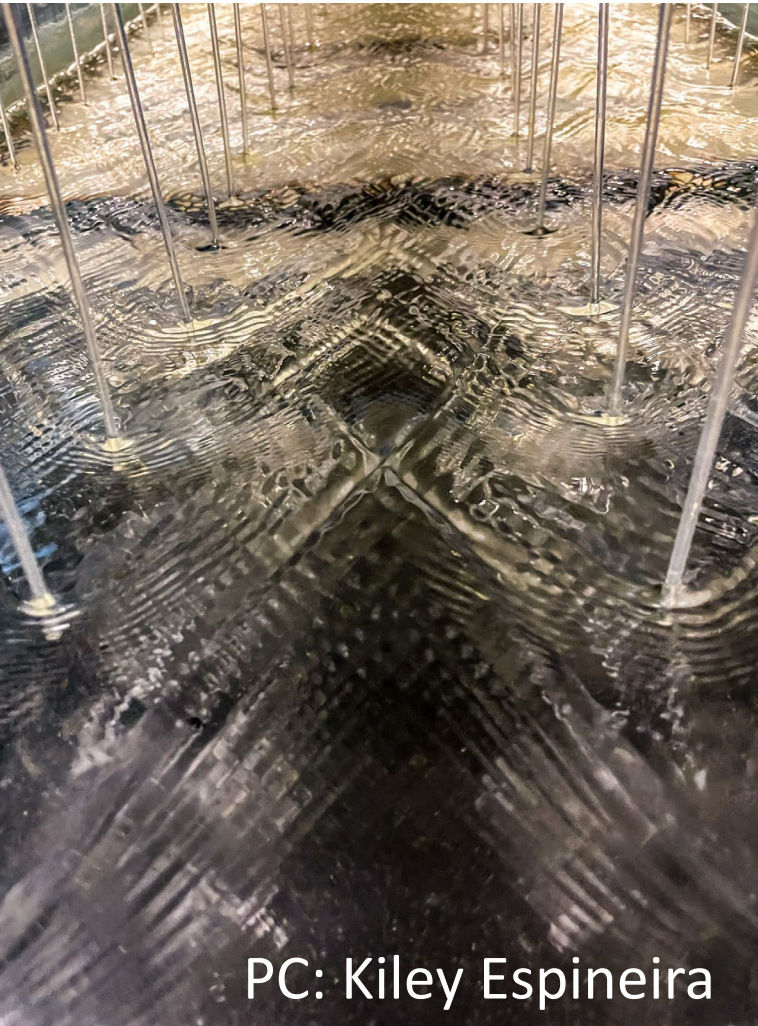
k is calculated from DO measurements



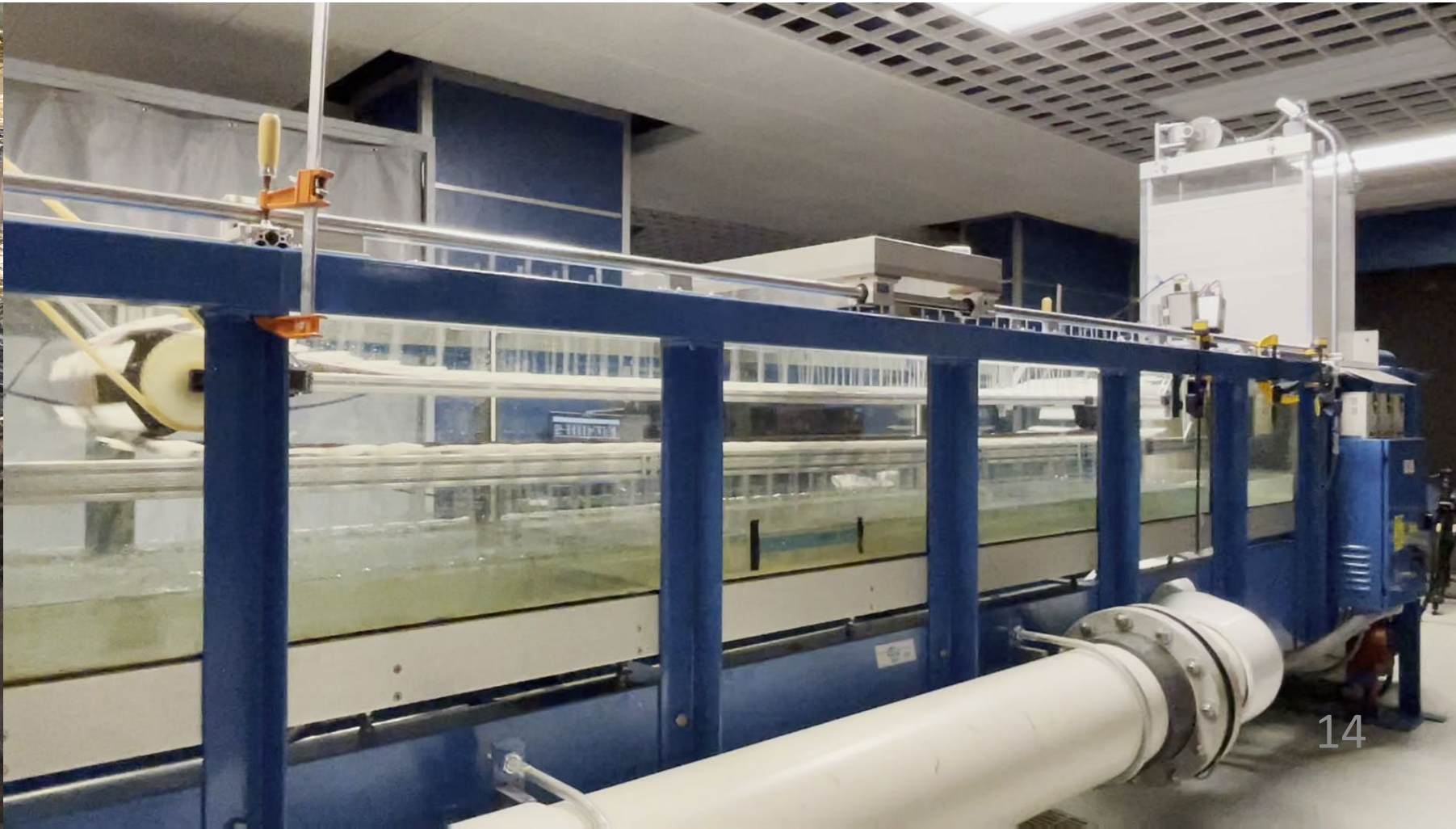
Results with static dowels: up to **66%** enhancement



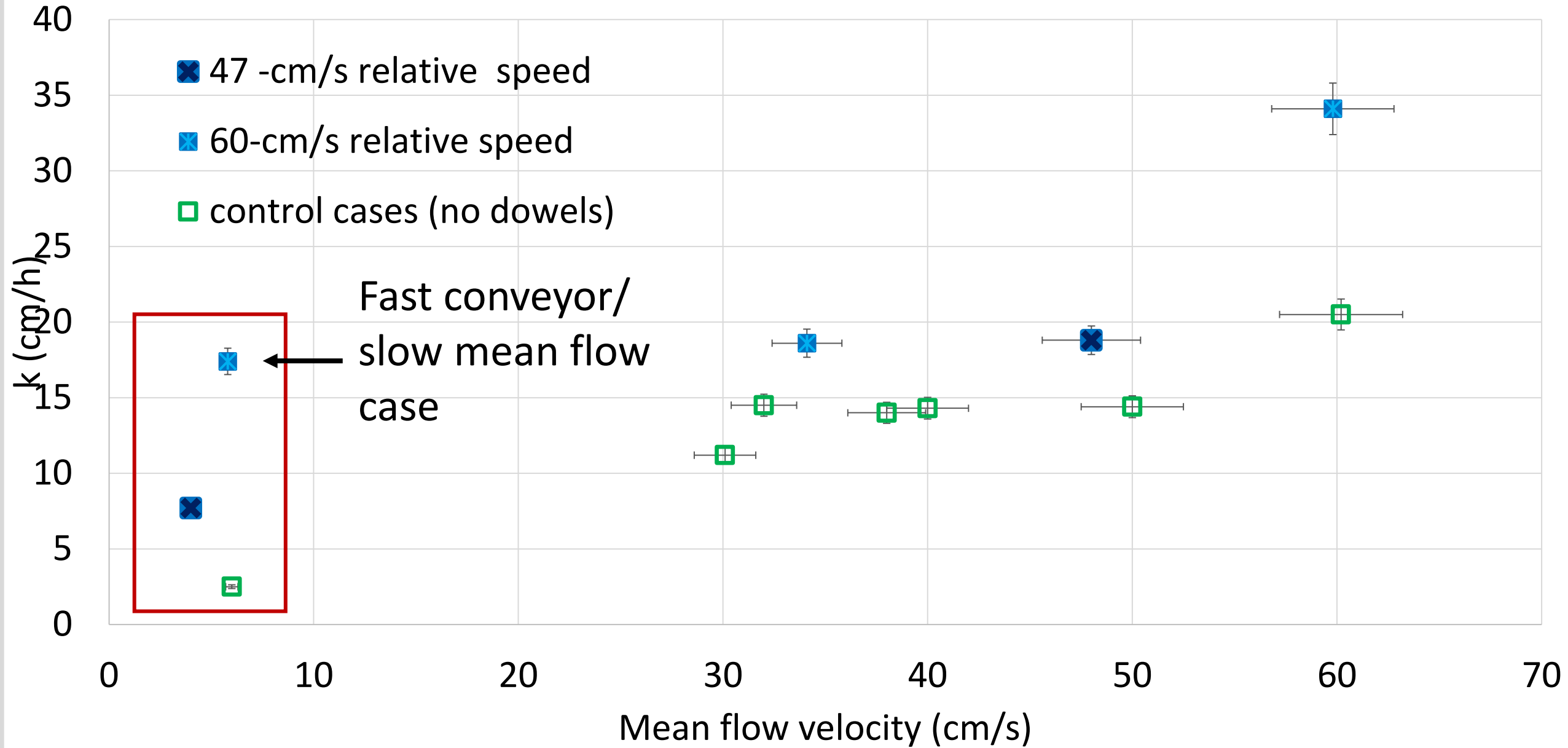
Conveyor belt isolates waves from background turbulence and BL dynamics



PC: Kiley Espineira



Moving dowels: Almost 6-fold increase in slow mean flow case

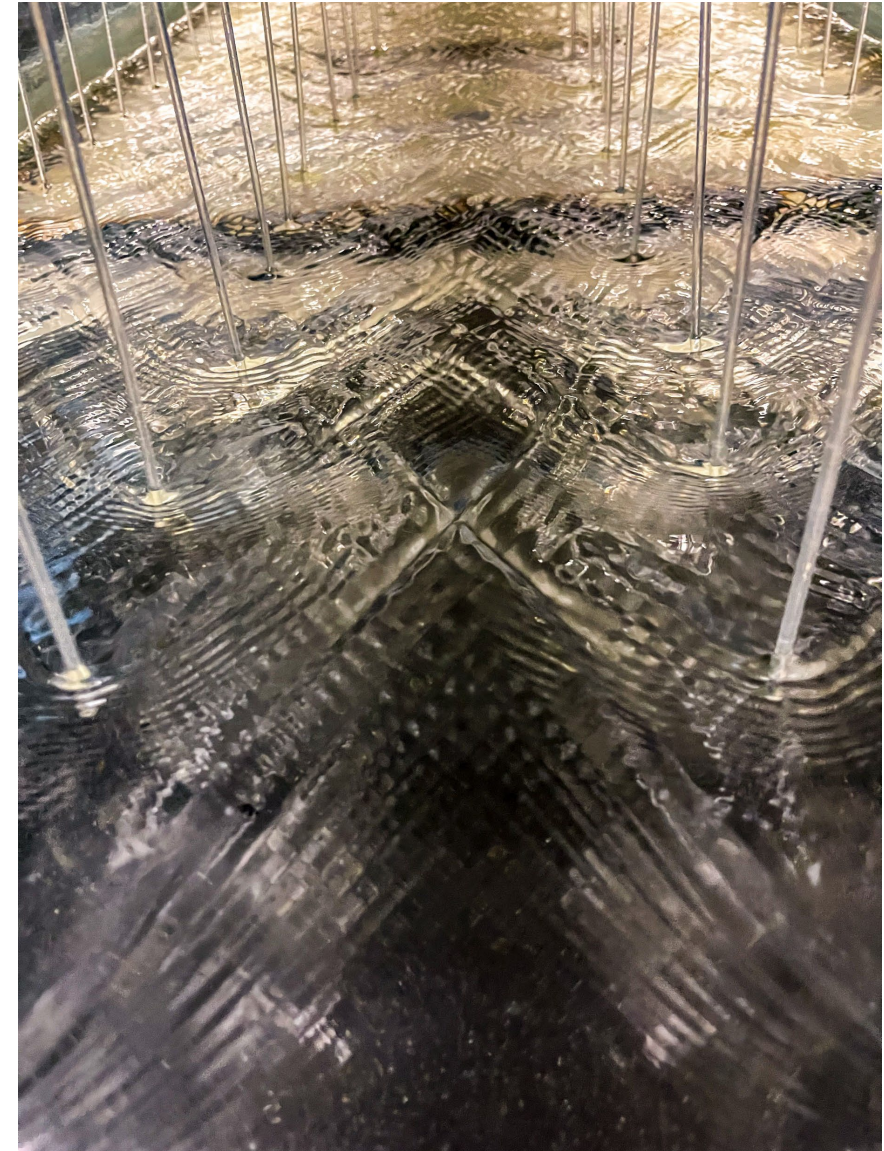


Summary

- Measured up to 67% enhancement in k due to static dowels at 60-cm/s flow speed.
- Measured a greater than 5-fold increase in k due to dowels conveyed at 60 cm/s against 5-cm/s flow speed.

Future Work

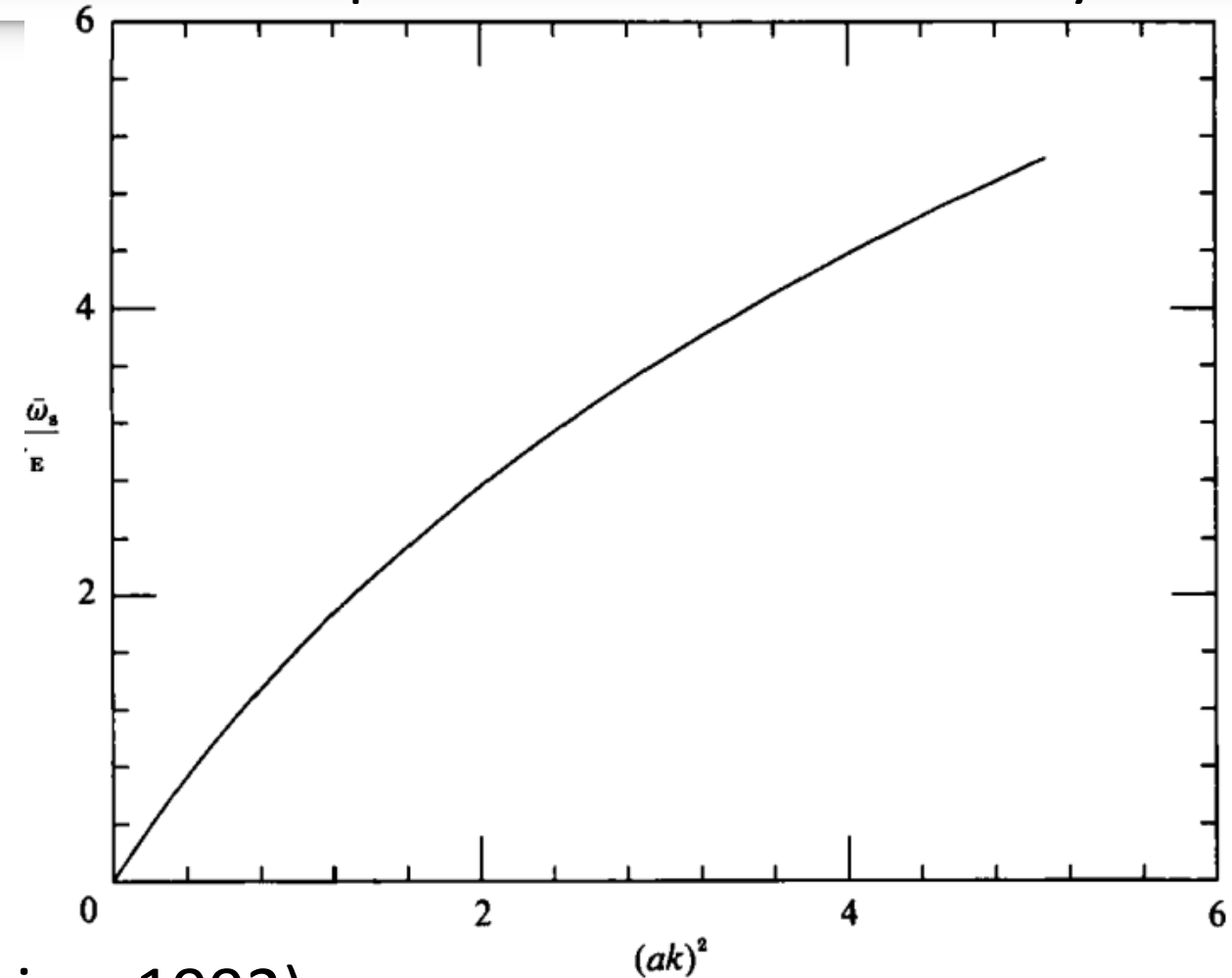
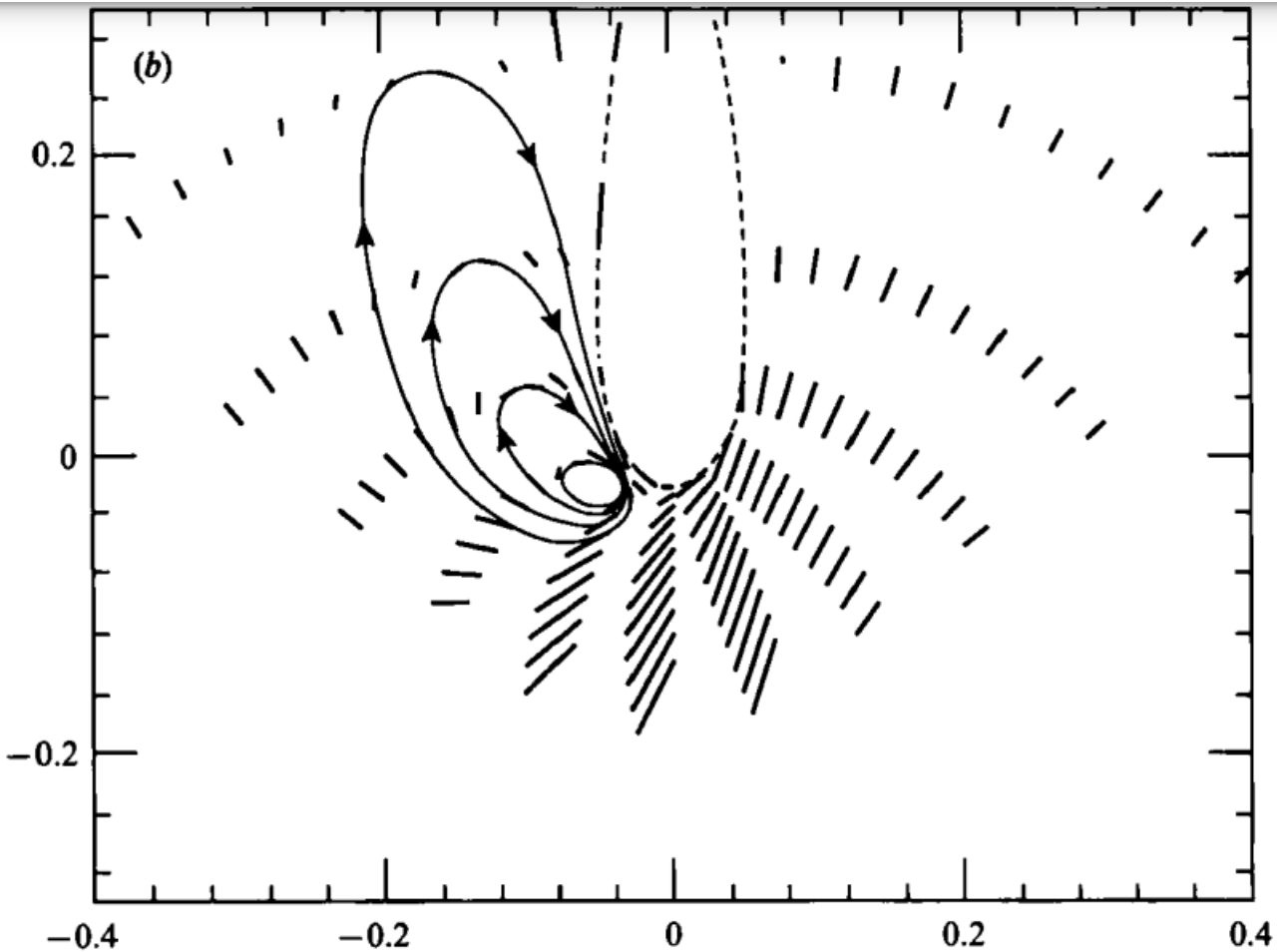
- Measure mean surface slope
- Use our dataset to develop a model that predicts these enhancements more broadly



Appendix

Capillary-gravity waves increase vorticity:

Steeper waves=more vorticity



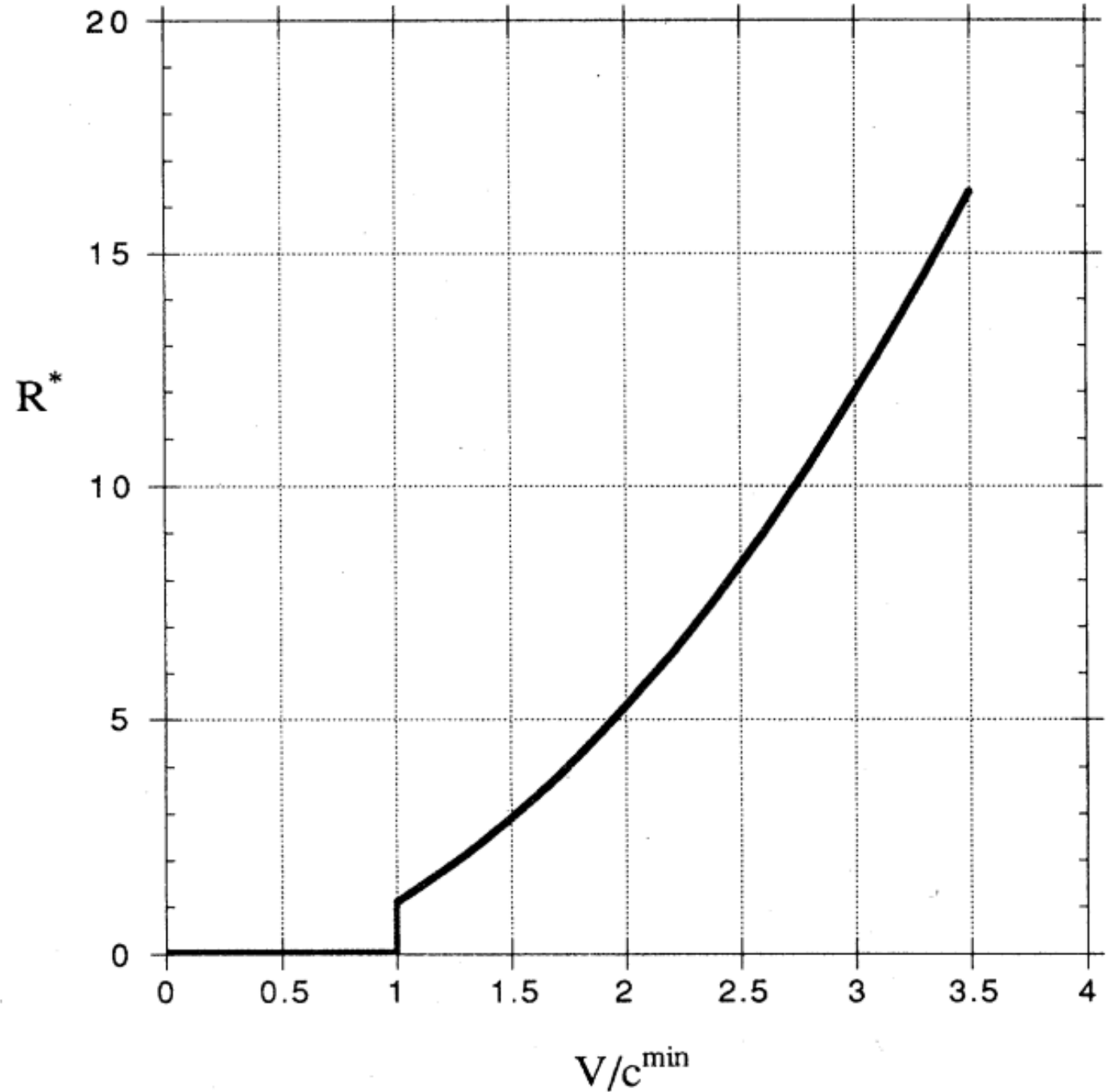
(Longuet-Higgins, 1992)

Moving dirac delta pressure field
(Raphael & De Gennes, 1996)

$$R^* = \pi \gamma (p^2 \kappa)^{-1} R$$

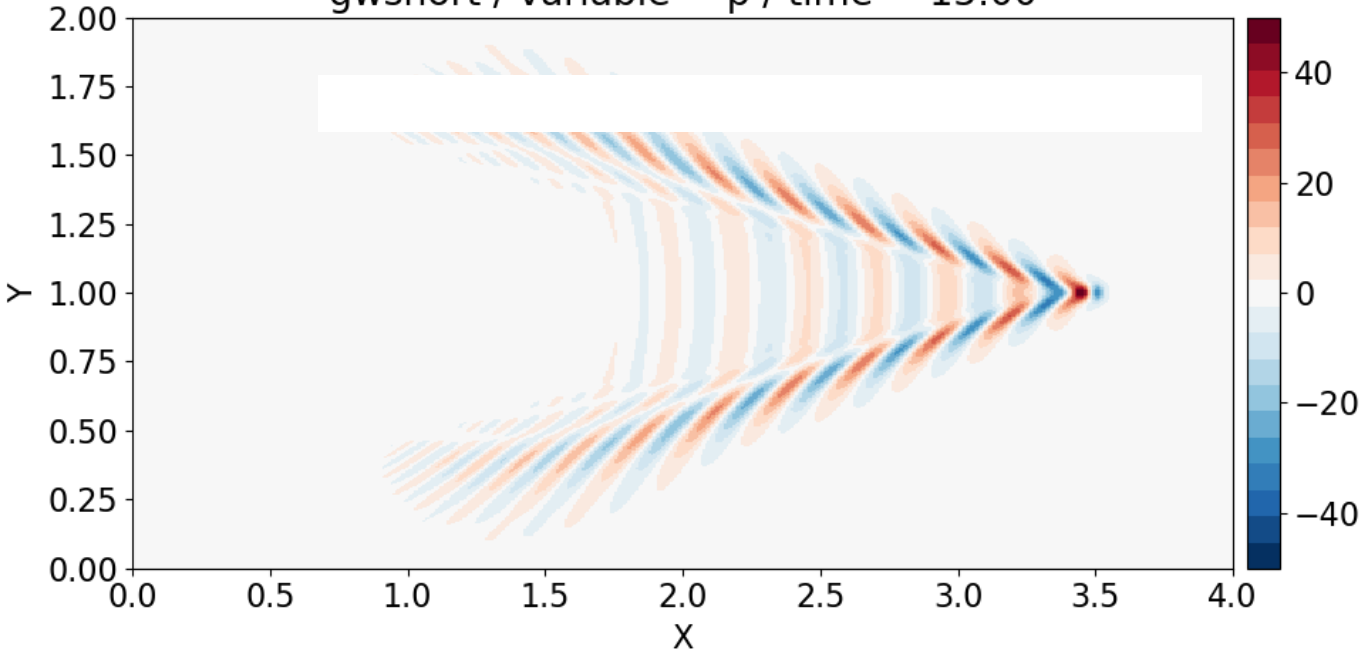
$$R = \frac{2}{3} \frac{p^2 \rho V^2}{\pi \gamma^2}, \quad V \gg c^{\min}$$

$$\kappa = (\rho g / \gamma)^{1/2}$$



Cannot analytically relate to amplitude yet

gwshort / variable = p / time = 15.00



(Chevy and Raphael, 2001)

Amplitude

$$A(\mathbf{K}_x, \mathbf{K}_y) = \frac{-ik_x V \hat{P}(\mathbf{K}_x, \mathbf{K}_y)}{\gamma k (\mathbf{K}^2 + \mathbf{K}_{min}^2) - \rho V^2 \mathbf{K}_x^2}$$

FT of Pressure field

$$\hat{P}(\mathbf{K}_x, \mathbf{K}_y) = F_0 \hat{\phi}(\mathbf{K}_x, \mathbf{K}_y)$$

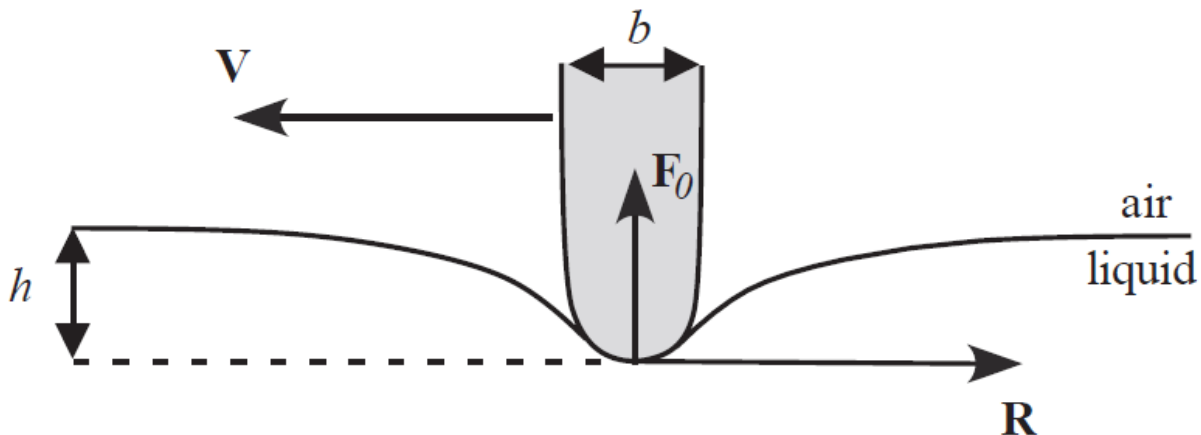
$F_0 = f(V)$

(in Rayleigh's linearized framework)

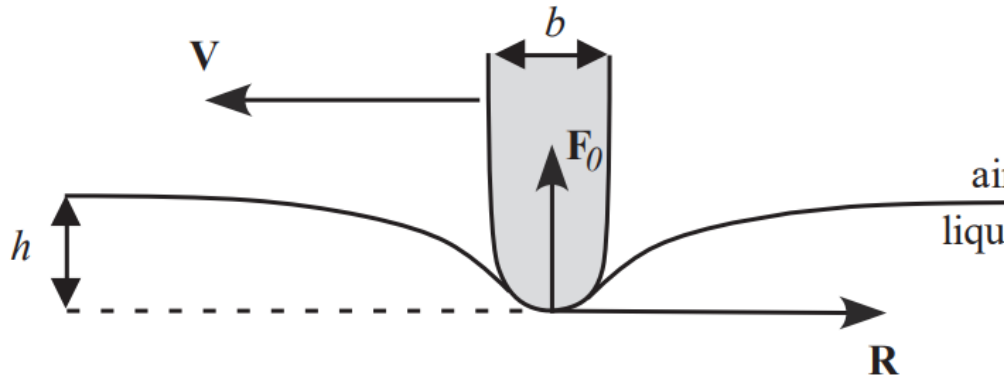
$$\hat{\xi}(\mathbf{K}_x, \mathbf{K}_y) = \frac{\frac{-F_0 K}{\rho} \hat{\phi}(\mathbf{K}_x, \mathbf{K}_y)}{\omega^2 - 4v^2 K^3 q + (2vK^2 - iV \cdot K)^2}$$

FT of free surface displacement

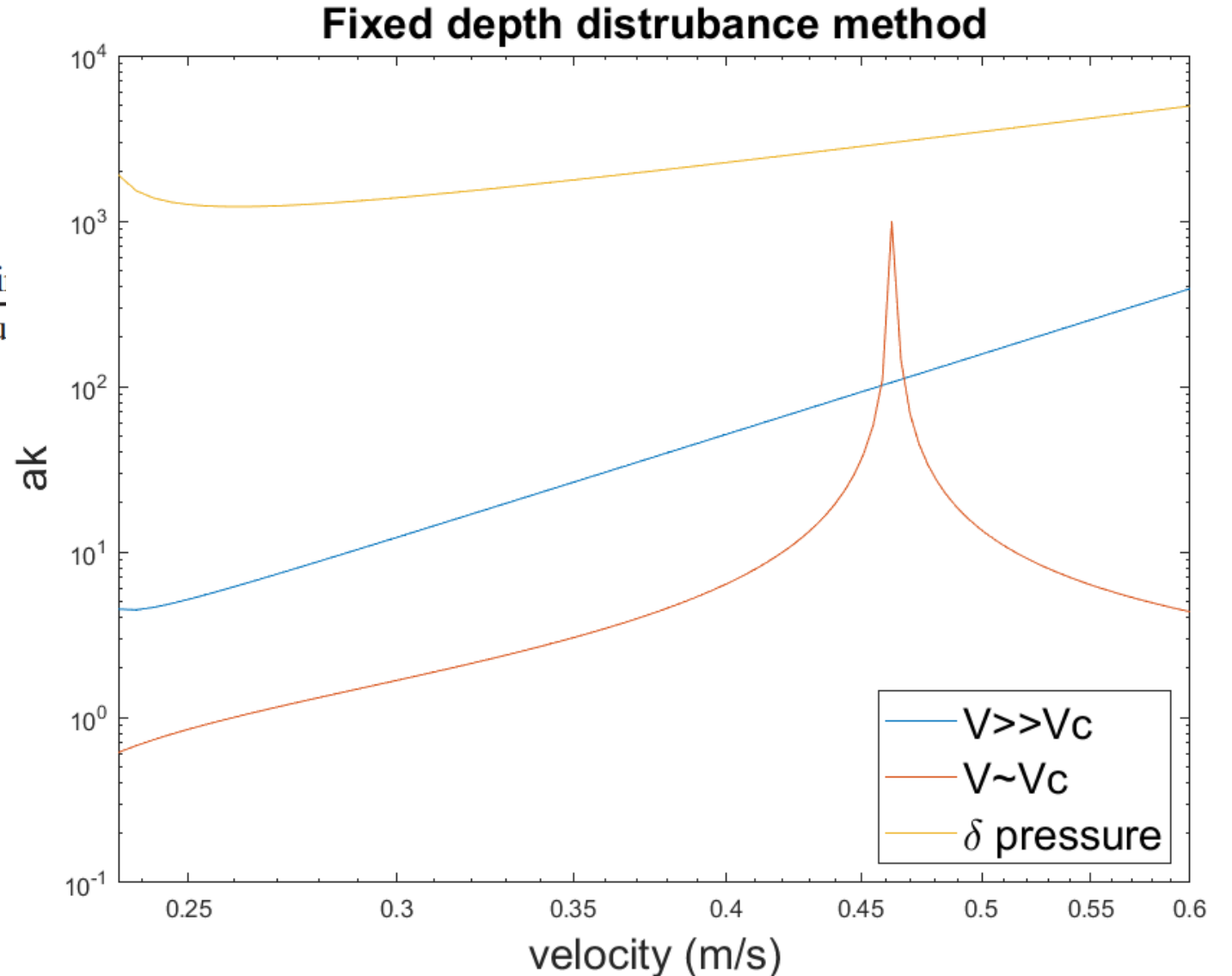
Dr. Guillaume Rouillet's wave2D model for gravity bow waves

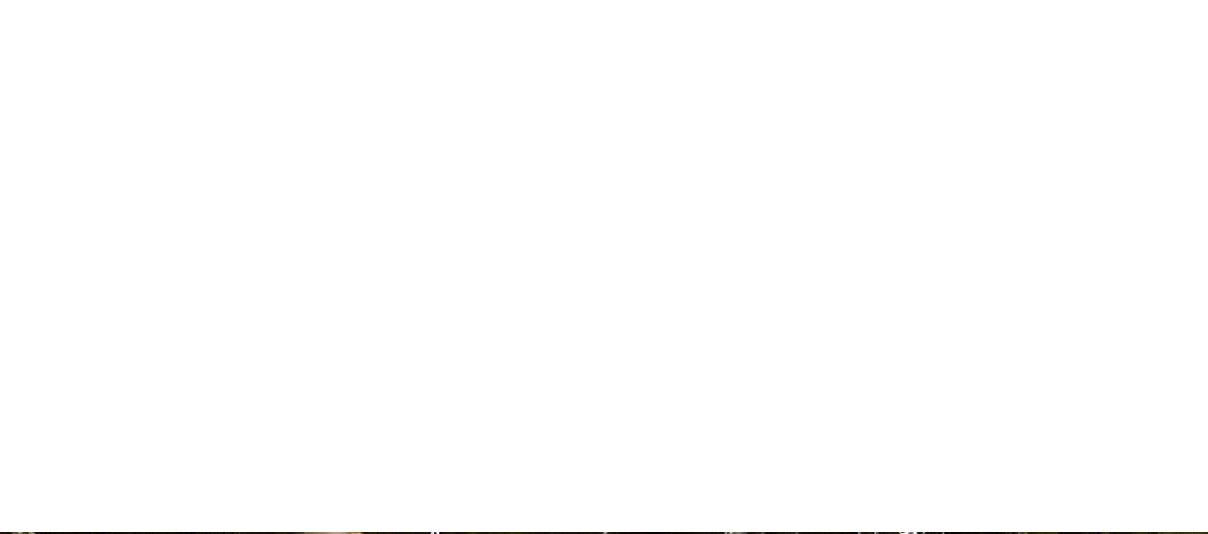


Relating steepness and flow velocity

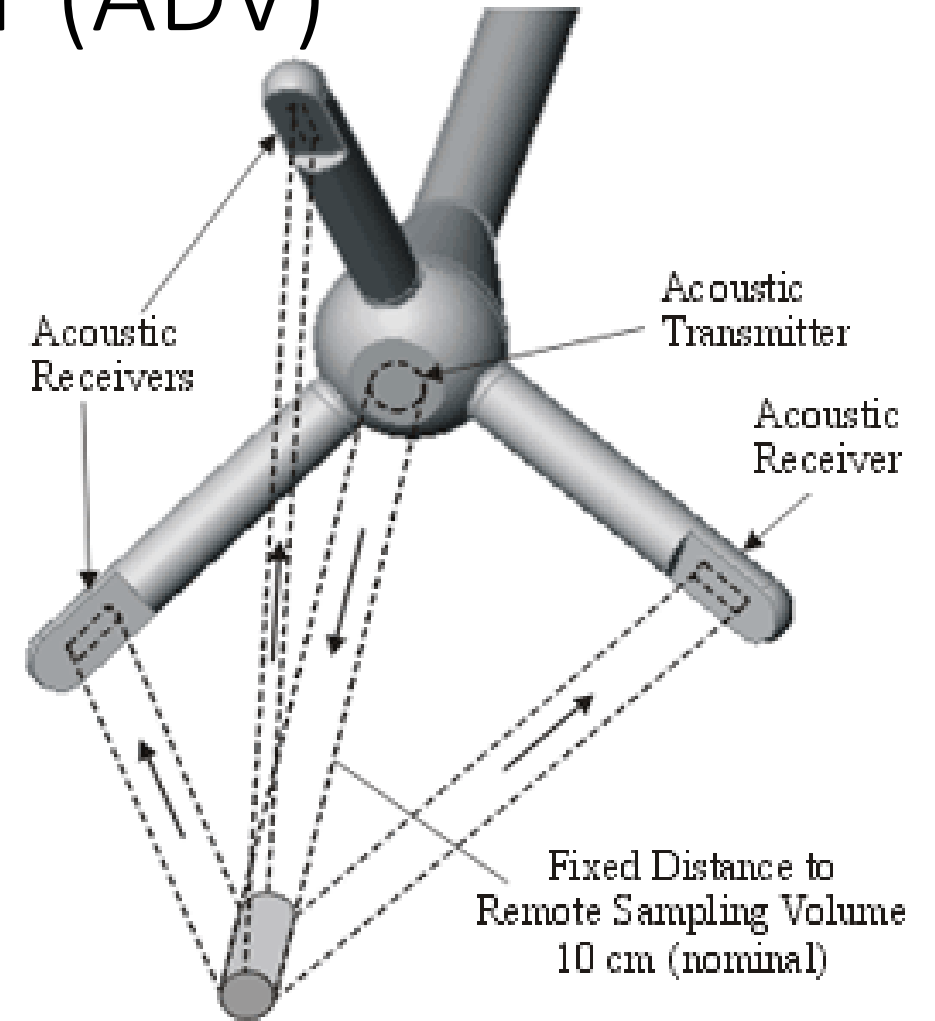
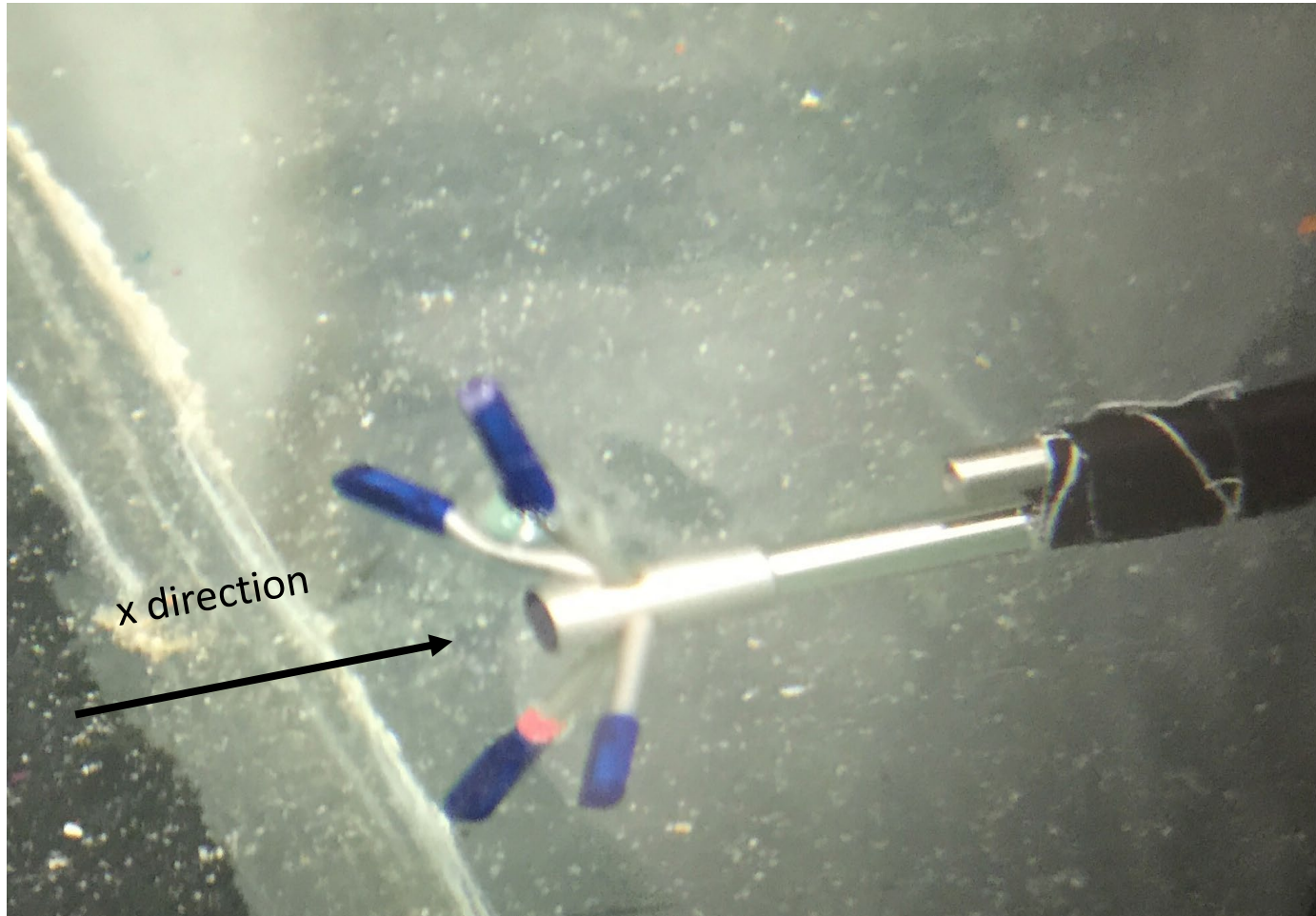


(Chevy & Raphaël, 2003;
Raphaël & de Gennes, 1996)





Acoustic Doppler Velocimeter (ADV)



Background: Surface Renewal Models

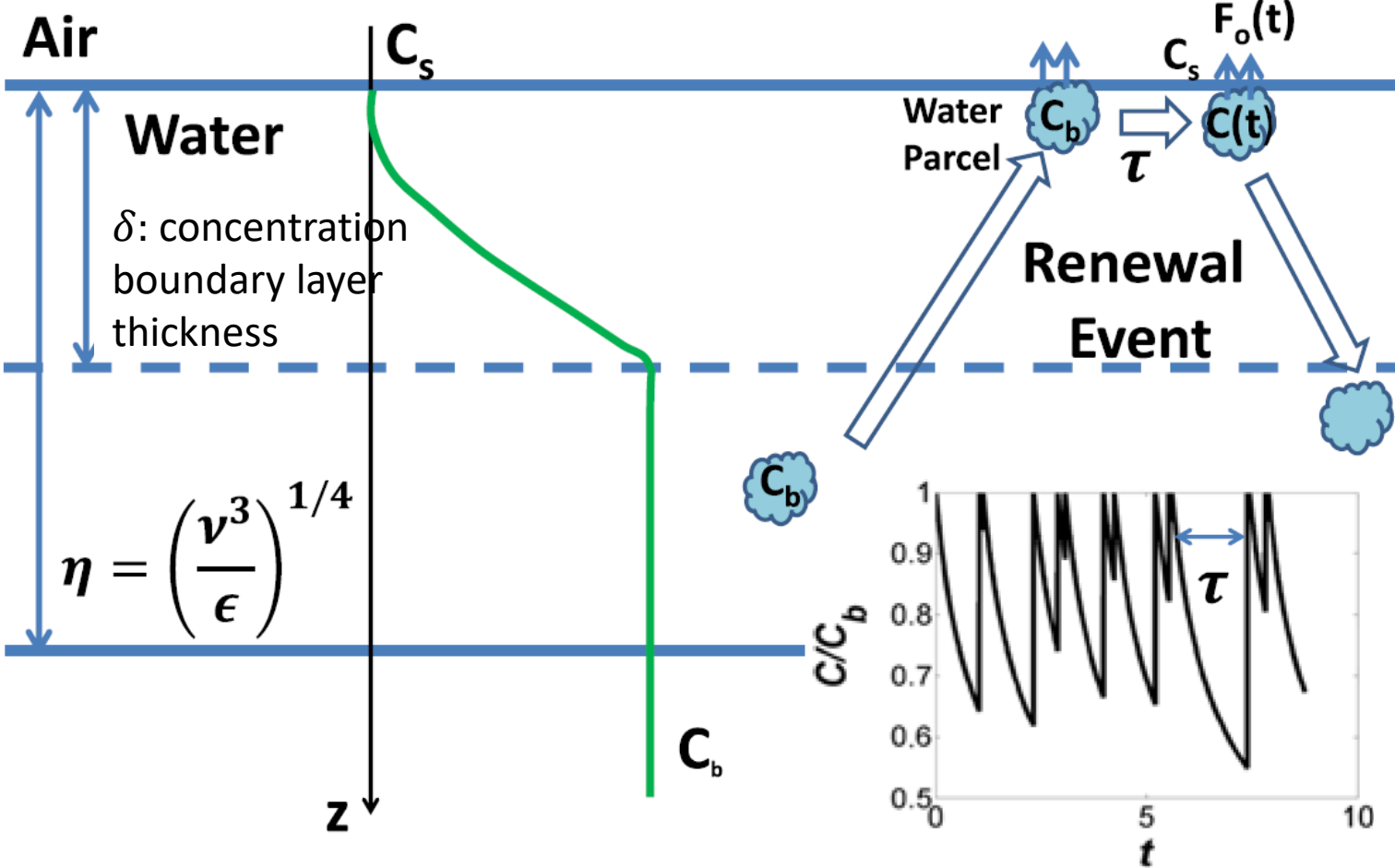
$$k \approx \sqrt{\frac{D}{\tau}}$$

τ is the average time between renewal events

For cases when the renewal time is much less than the molecular diffusion time scale:

$$\tau \ll T_D$$

$$T_D = \frac{\delta^2}{D}$$



(Katul and Liu, 2017)

Background: Surface Renewal Models

(Danckwerts, 1951)

Large Eddy Model:

$$\tau = \frac{L}{\sqrt{u'^2}}$$

where L is the turbulence integral length scale and $\sqrt{u'^2}$ is the turbulence intensity or standard deviation of velocity, u (Fortescue and Pearson, 1967)

Use when $Re_t = \frac{\sqrt{u'^2}L}{\nu} < 500$ (Theofanous *et al.*, 1975)

Small Eddy Model: (Lamont and Scott, 1970)

$$\tau = \sqrt{\frac{\nu}{\varepsilon}} = \eta_T,$$

where ν is the kinematic viscosity and ε is the energy dissipation rate (Banerjee *et al.*, 1968 and Lamont and Scott, 1970)

Use when $Re_t = \frac{\sqrt{u'^2}L}{\nu} > 500$ (Theofanous *et al.*, 1975)

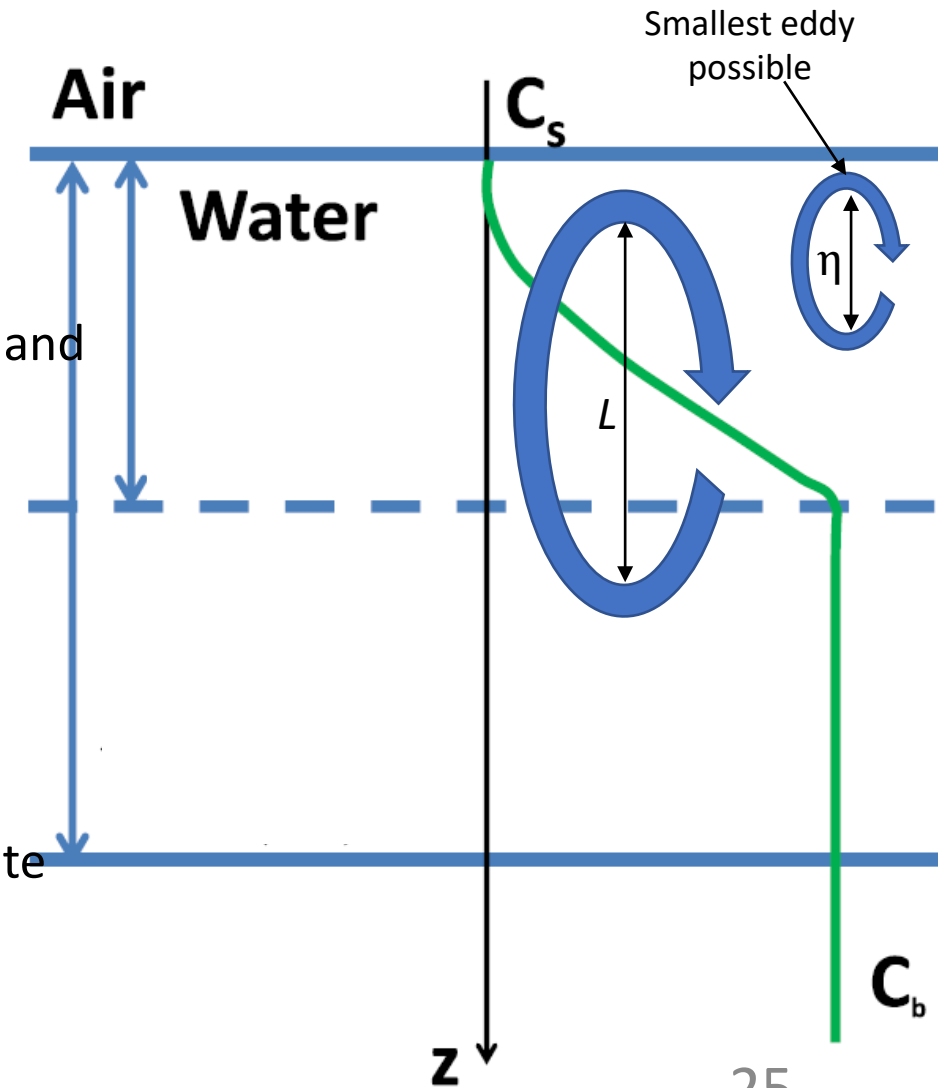
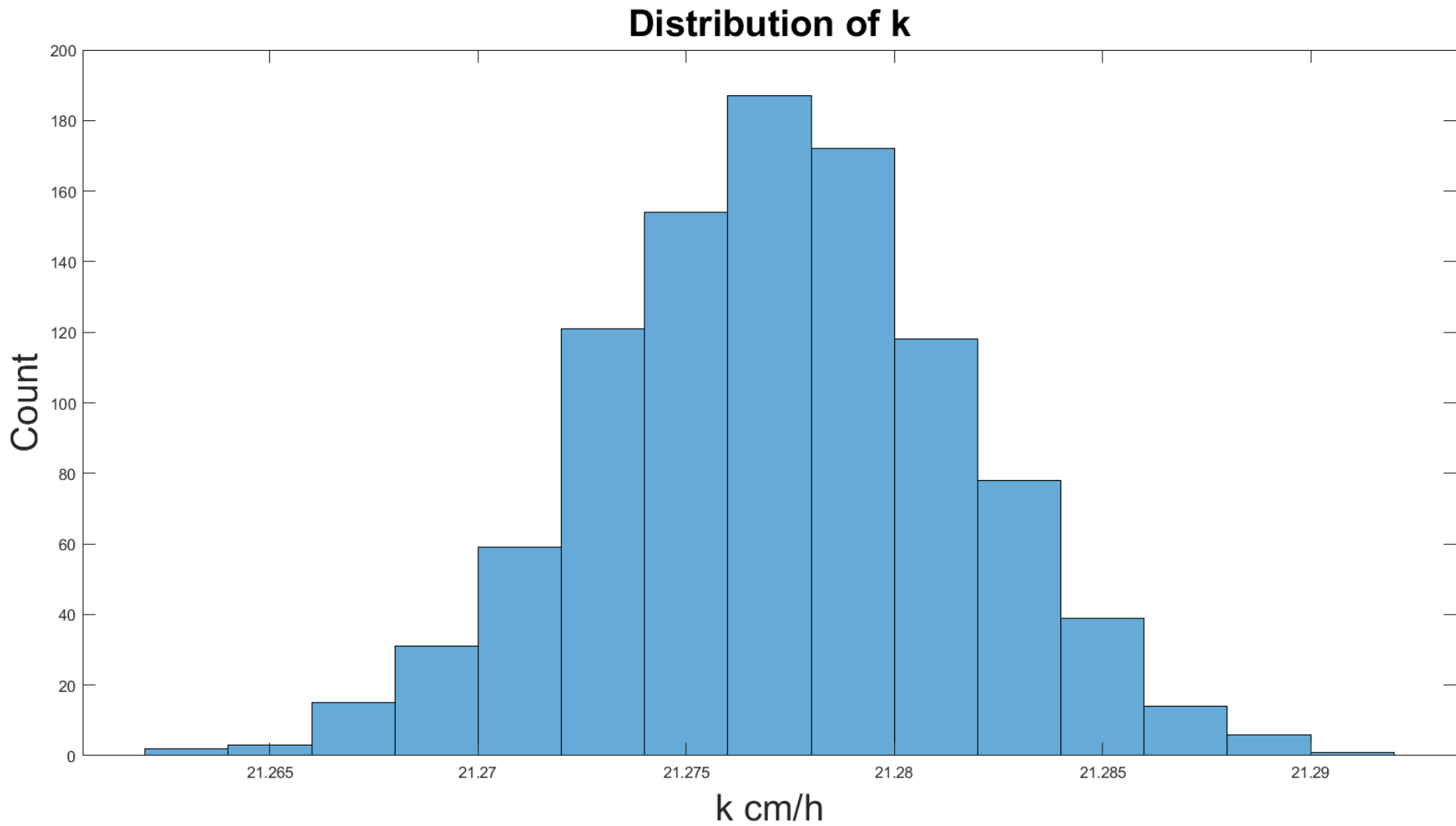


Figure modified from Katul and Liu (2017). Eddies not to scale



Resampling 1,000 times yields:

$$\bar{k} = 21.2772 \frac{cm}{h}, \quad \sigma_k = 0.0044 \frac{cm}{h}, \quad k(975) = 21.2857 \frac{cm}{h}, \quad k(25) = 21.2684 \frac{cm}{h}$$

$$k = 21.277 \pm 0.009 \frac{cm}{h}$$