## Air-sea scalar transfer – effects of wind and waves on equivalent roughness length

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# Parameterization of air-sea fluxes of momentum and scalar (temperature, humidity, soluble gases)

1. We parameterize the fluxes using mean wind speed at 10 meter height  $(u_{10})$ , mean scalar concentrations at 10 meter height  $(\theta_{10})$  and at sea surface  $(\theta_0)$ :

$$u_*^2 = C_d u_{10}^2, \qquad \theta_* u_* = C_{\theta} (\theta_{10} - \theta_0) u_{10}$$

2. We assume logarithmic profiles (in height z) of mean wind speed (u) and mean scalar concentration ( $\theta$ ), under neutral conditions (air temperature = water temperature), inside the constant stress (flux) layer, but above the direct effects of surface waves:

$$\frac{u(z)}{u_*} = \frac{1}{\kappa} \log \frac{z}{z_0} , \qquad \frac{\theta(z) - \theta_0}{\theta_*} = \frac{0.74}{\kappa} \log \frac{z}{z_0}$$

Then,  $C_d$  and  $C_\theta$  are expressed as:

$$C_{d} = \left(\frac{1}{\kappa}\log\frac{10}{z_{0}}\right)^{-2}, \qquad C_{\theta} = \left(\frac{1}{\kappa}\log\frac{10}{z_{0}}\right)^{-1} \left(\frac{0.74}{\kappa}\log\frac{10}{z_{\theta}}\right)^{-1}$$

 $C_d$  is a function of  $z_0$  alone, but  $C_{\theta}$  is a function of both  $z_0$  and  $z_{\theta}$ .

 $\tau = \rho_a u_*^2$ : wind stress  $\rho_a$ : air density  $u_*$ : friction velocity  $\theta_* u_*$ : scalar flux

 $C_d$ : drag coefficient  $C_{\theta}$ : scalar transfer coefficient

 $z_0$ : equivalent surface roughness of wind  $z_{\theta}$ : equivalent surface roughness of scalar We know how  $C_d$  and  $C_{\theta}$  vary with wind speed  $u_{10}$  (on average) quite well, except at very high wind speeds where the uncertainty is large.



Up to wind speed 20 m/s,  $C_d$  rapidly increases but  $C_{\theta}$  remains almost constant.

Recall that  $C_d$  and  $C_{\theta}$  are expressed as:

$$C_{d} = \left(\frac{1}{\kappa}\log\frac{10}{z_{0}}\right)^{-2}, \qquad C_{\theta} = \left(\frac{1}{\kappa}\log\frac{10}{z_{0}}\right)^{-1} \left(\frac{0.74}{\kappa}\log\frac{10}{z_{\theta}}\right)^{-1}$$

- From COARE3.5  $C_d$  and COARE3.0  $C_{\theta}$ , we may calculate both  $z_0$  and  $z_{\theta}$  as functions of  $u_{10}$ .
- $z_0$  rapidly increases with  $u_{10}$ .
- $z_{\theta}$  rapidly decreases with  $u_{10}$ .
- I will demonstrate how ocean surface waves are responsible for increasing z<sub>0</sub> and decreasing z<sub>θ</sub>.



 $z_0$  (blue) and  $z_{\theta}$  (red) as functions of  $u_{10}$ 

## Sea state dependence of $C_d$

- If ocean surface waves are responsible for increasing z<sub>0</sub>, it is expected that different sea states yield different z<sub>0</sub> (and C<sub>d</sub>), even if wind speed u<sub>10</sub> is the same.
- However, the dependence of z<sub>0</sub> (and C<sub>d</sub>) on wave parameters (e.g., wave age, wave steepness) is not always clear and not well constrained.



Wave age dependence of normalized roughness  $z_0$  (Edson et al. 2013)

## Sea state dependence of $C_d$

Our recent (observation + model) study suggests  $C_d$  weakly depends on  $u_{10}$  but strongly depends on sea states under tropical cyclones (wind speed  $u_{10}$  25-55 m/s) (Zhou et al. 2022).



between wind and dominant waves

## Sea state dependence of $C_d$

Our recent modeling study suggests  $C_d$  strongly depends on sea states in coastal shallow waters (Chen et al. 2020a,b)

Right panels show that  $C_d$  under steady uniform wind significantly increases in shallow waters.

Bottom panels show how sea state dependence of  $C_d$  under tropical cyclones is enhanced in shallow waters.





Sea state dependence of  $C_{\theta}$  and  $z_{\theta}$ ?

$$C_{\theta} = \left(\frac{1}{\kappa}\log\frac{10}{z_{0}}\right)^{-1} \left(\frac{0.74}{\kappa}\log\frac{10}{z_{\theta}}\right)^{-1}$$

Two key questions:

- If  $C_d$  and  $z_0$  are sea state dependent, are  $C_{\theta}$  and  $z_{\theta}$  also sea state dependent?
- Can the sea state dependence of  $z_{\theta}$  compensate (cancel) the sea state dependence of  $z_0$ , so that  $C_{\theta}$  remains sea state independent?

## Large eddy simulation (LES) of wind over a sinusoidal wave train

Surface wave length  $\lambda$ , wavenumber  $k = \frac{2\pi}{\lambda}$ , wave amplitude a, wave phase speed cWave slope ka = 0.27

LES domain size:  $5\lambda \times 5\lambda \times 2.4\lambda$ , LES grids:  $256 \times 256 \times 256$ 

Wind is driven by externally imposed horizontal pressure gradient

Weak heating from bottom for scalar flux simulations

Background surface roughness (due to unresolved small waves)  $z_{0b}$ ,  $kz_{0b} = 2.7 \times 10^{-3}$ 

Background surface roughness of scalar (due to unresolved small waves)  $z_{\theta b}$ We vary (relative) wind forcing ( $c/u_*$ ), wind direction, and  $z_{\theta b}$ 

#### **REFERNECES:**

LES studies: Hara and Sullivan 2015; Sullivan et al. 2017; Husain et al. 2022a; 2022b Validation of LES against laboratory observation: Husain et al. 2019

### Large eddy simulation (LES) of wind over a sinusoidal wave train

Phase averaged flow fields,  $c/u_*=1.4$  (strongly wind forced).

WIND from Left to Right →



Behind the crest is a region of reduced wind speed, reduced turbulence, and reduced scalar concentration gradient.

#### **Airflow separation**

Although we simulate conditions where the "ensemble averaged" flow field does not separate over a non-breaking wave train, instantaneous flow fields often show "separation like" patterns.

Our LES results are consistent with laboratory observations by Buckley and Veron (Univ. of Delaware).



Top: PIV observations, Bottom: LES results

### Large eddy simulation (LES) of wind over a sinusoidal wave train Phase averaged flow fields



#### Airflow frequently separates from the crest.

Behind the crest is a region of reduced wind speed, reduced turbulence, and reduced scalar concentration gradient.

### Large eddy simulation (LES) of wind over a sinusoidal wave train Phase averaged flow fields



Airflow frequently separates from the crest.

Behind the crest is a region of reduced wind speed, reduced turbulence, and reduced scalar concentration gradient, where the surface tangential stress becomes close to zero.

## Horizontally averaged analyses using a mapped vertical coordinate $\zeta$ (Hara and Sullivan, 2015).

How are the mean wind speed profile  $\langle u(\zeta) \rangle$  modified from a flat surface to a wavy surface (if wind stress and scalar flux remain the same)?



 $\langle u \rangle$  over a wave train (blue) is shifted to the left of  $\langle u \rangle$  over a flat surface (red) away from the surface.

Roughness  $z_0$  increases over a wave train.

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Roughness  $z_0$  increases over a wave train.

Wind shear is reduced over a wave train (blue) than over a flat surface (red) near the surface, hence,  $z_0$  increases.

Wind shear is enhanced near the height of wave crest. (Airflow separation makes the surface smoother?)

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How are the mean wind speed profile  $\langle u(\zeta) \rangle$  modified from a flat surface to a wavy surface (if wind stress and scalar flux remain the same)?



Momentum flux budget Green: wave coherent stress Blue: turbulent stress Red: total stress

Turbulent stress is significantly reduced near the surface, causing the reduced wind shear.



Wind shear is reduced over a wave train (blue) than over a flat surface (red) near the surface, hence,  $z_0$  increases.

Wind shear is enhanced near the height of wave crest. (Airflow separation makes the surface smoother?)

We may estimate  $z_0/z_{0b}$ , i.e., how much roughness has increased over a wave train compared to a flat surface. (Husain et al. 2022a,b)



With waves following wind (black), roughness is more enhanced with stronger wind forcing.

With wave opposing wind (red), roughness is more enhanced with weaker wind forcing.



wind and wave misaligned ightarrow

If wave direction is close to wind direction, roughness is enhanced.

If wave direction is misaligned from wind direction by more than 45 degrees, roughness is reduced compared to a flat surface.



 $\langle u \rangle$  over a wave train (blue) is shifted to the left of  $\langle u \rangle$  over a flat surface (red) away from the surface.

Roughness  $z_0$  has increased by a factor of ~5.  $C_d$  has increased by ~40%  $\langle \theta \rangle$  over a wave train (blue) is shifted to the right of  $\langle \theta \rangle$  over a flat surface (red) away from the surface.

Roughness  $z_{\theta}$  has decreased by a factor of ~5.  $C_{\theta}$  remains almost unchanged.



 $\langle \theta \rangle$  is shifted to the left because the vertical gradient of  $\langle \theta \rangle$  is significantly enhanced near the surface.

 $\langle \theta \rangle$  over a wave train (blue) is shifted to the left of  $\langle \theta \rangle$  over a flat surface (red) away from the surface.

Roughness  $z_{\theta}$  has decreased.



 $\langle \theta \rangle$  is shifted to the left because the vertical gradient of  $\langle \theta \rangle$  is significantly enhanced near the surface.

Scalar flux budget Green: wave coherent flux Blue: turbulent flux Red: total flux

Turbulent scalar flux is not reduced near the surface. But, turbulent stress is reduced near the surface. Therefore, the vertical gradient of  $\langle \theta \rangle$  increases. (Near the surface,  $u_*$  decreases but  $\theta_* u_*$  does not change. Therefore,  $\theta_*$  increases.)



 $\langle u \rangle$  over a wave train (blue) is shifted to the left of  $\langle u \rangle$  over a flat surface (red) away from the surface.

 $\langle \theta \rangle$  over a wave train (blue) is shifted to the right of  $\langle \theta \rangle$  over a flat surface (red) away from the surface.

Roughness  $z_0$  has increased.

Roughness  $z_{\theta}$  has decreased.

Increase of  $z_0$  and decrease of  $z_{\theta}$  seem to (roughly) cancel each other, i.e.,  $C_{\theta}$  remains (roughly) unchanged, if the background roughness  $z_{0b}$  and  $z_{\theta b}$  are the same.

Recall that  $z_0$  is much larger than  $z_{\theta}$  except for very low wind speeds.

Therefore, the background roughness (due to unresolved small waves)  $z_{\theta b}$ should be much less than the background roughness  $z_{0b}$ .



 $z_0$  (blue) and  $z_{\theta}$  (red) as functions of  $u_{10}$ 

We repeat the same LES experiment but reduces  $z_{\theta b}$  ( $z_{\theta b}/z_{0b} = 0.01$ ), which is more realistic.



Increase of  $z_0$  and decrease of  $z_{\theta}$  do not cancel each other, and  $C_{\theta}$  is significantly (~20%) reduced by the waves, if the background roughness  $z_{\theta b}$  is much less than the background roughness  $z_{0b}$ .

### Conclusion:

- Scalar transfer coefficient  $C_{\theta}$  is a function of both  $z_0$  and  $z_{\theta}$ .
- Recent studies show that  $z_0$  can be strongly sea state dependent, particularly under tropical cyclones and in coastal shallow waters.
- Our LES study shows that  $z_0$  is significantly modified (usually increased) over a sinusoidal wave train (compared to a flat surface), depending on wind strength and direction. Airflow separations may play a significant role in modifying  $z_0$ .
- Our preliminary LES study shows that:
  - When  $z_0$  increases over a sinusoidal wave tarin (because the turbulent stress is reduced),  $z_{\theta}$  decreases (because the scalar flux is NOT reduced).
  - The effects of increasing  $z_0$  and decreasing  $z_{\theta}$  on  $C_{\theta}$  do not always cancel.  $C_{\theta}$  can be significantly reduced by waves.